

5.4 Alternative ultra-light DM candidates

5.4.1 Axion-like particles

Goldstone theorem:

A spontaneously broken global symmetry gives rise to a massless "Goldstone" boson a

\Rightarrow Shift symmetry $a \rightarrow a + \epsilon$

If the global symmetry is not exact, the new boson acquires a small mass

("Pseudo-Goldstone" boson)

Examples: - Axion

Global symmetry: Peccei-Quinn sym
broken at f_a

Explicit breaking: QCD effects at
scale f_π

$$\Rightarrow m_a \sim \frac{f_\pi}{f_a} \ll 1$$

- Majorons

Global symmetry: Lepton number

\hookrightarrow Spontaneously broken in
models of neutrino mass
generation

- String theory axions

String theory requires 10 dim.
out of which 6 need to
be compactified



Closed cycles give rise to spin-0
fields with vanishing potential.

Realistic string theories can have
100s of axions ("Axiverse")

Allowed interactions of PGBs
dictated by shift symmetry

$$\mathcal{L}_{\text{int}} = \frac{1}{32\pi^2} \frac{a}{f_a} G^{a\mu\nu} \tilde{G}_{\mu\nu}^a - \frac{1}{4} g_{\text{af}}^2 F^{\mu\nu} \tilde{F}_{\mu\nu} \\ + \sum_e g_{ae} \partial_n^a \bar{e} \gamma^\mu \gamma^5 e + \sum_q g_{aq} \partial_n^a \bar{q} \gamma^\mu \gamma^5 q$$

- g_{af} , g_{ae} , g_{aq} are typically of order $\frac{1}{f_a}$
but vary from model to model

- m_a is a free parameter!

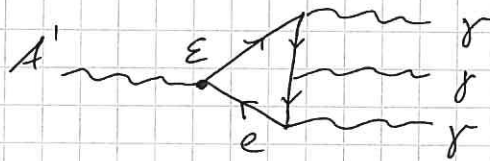
5.4.2 Hidden photons

Similar to the dark photon A' discussed in §4.2 + 4.3

\Rightarrow New $U(1)'$ gauge boson coupling to SM particles via kinetic mixing ϵ

\Rightarrow Very small mass $m_{A'} \ll 1 \text{ MeV}$

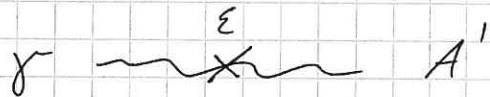
\hookrightarrow Only allowed decay channel



\Rightarrow Lifetime long compared to age of Universe

Production in the early Universe via non-thermal processes:

oscillations



coupling to electrons



Can be all of DM for $\epsilon \sim 10^{-12} - 10^{-9}$

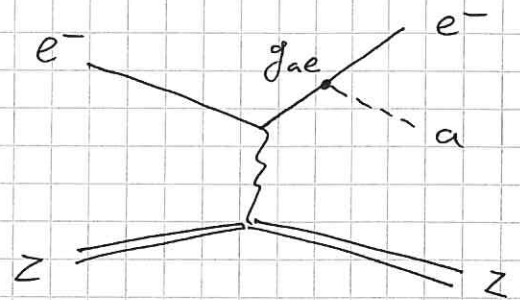
5.5 Astrophysical constraints on axions (and other ultra-light particles)

Stars can produce axions via a number of processes:

Primakoff process:



Bremsstrahlung:



Compton scattering:



Typical energy \sim core temperature of star

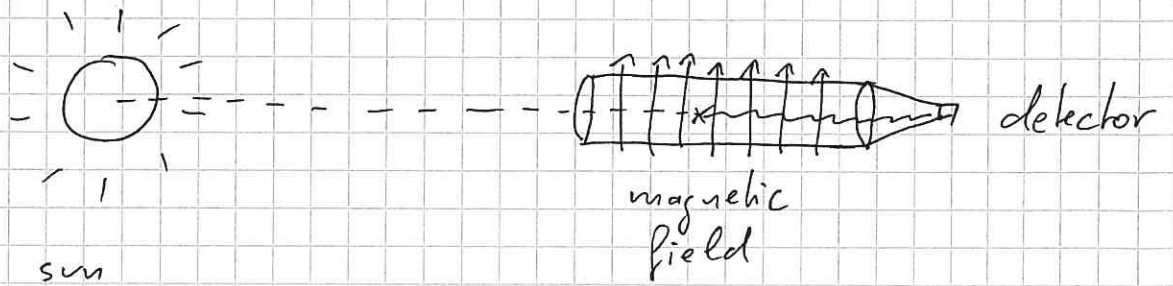
$$E_a \sim T_{\text{star}} \quad (\text{e.g. } T_{\text{sun}} \sim 1 \text{ keV})$$

For $g_{\gamma a} = 10^{-10} \text{ GeV}^{-1}$ the axion flux corresponds to 0.2% of sun's luminosity

$$\Rightarrow \text{Flux on Earth: } \Phi_a \sim 4 \cdot 10^{11} \text{ cm}^{-2} \text{ s}^{-1}$$

5.5.1 Helioscopes

Basic idea:



As for LSW experiments,

$$P_{\text{sig}} \sim \mathcal{F}(q \cdot L) \quad \text{with} \quad q = \frac{m_a^2}{2E_a}$$

which requires $q \cdot L \lesssim 1$

But now $E_a \sim \text{keV}$ (and $L \sim 10 \text{ m}$)

\Rightarrow Sensitivity up to $m_a \sim 10 \text{ meV}$

\hookrightarrow Even larger masses possible
when using buffer gas

\Rightarrow Can probe heavier axions
than other searches

CAST (at CERN)

- $B \sim 9 \text{ T}$

- Trace sun for several hours per day

- Constraint: $g_{\text{a}\gamma} \lesssim 10^{-10} \text{ GeV}^{-1}$

IA XO will improve this to $5 \cdot 10^{-12} \text{ GeV}^{-1}$

(98) \Rightarrow probe QCD axions in meV range

5.5.2 Axioelectric effect

If $g_{ae} \neq 0$, axions can ionize neutral atoms



⇒ Search for ionization events in low-background detectors

↳ Same technology as for WIMP direct detection

LUX: $g_{ae} \lesssim 10^{-8} \text{ GeV}^{-1}$

5.5.3 Horizontal Branch (HB) stars

All stars can produce axions

Even if the axions are undetected, they provide an additional energy loss mechanism

⇒ Stars cool faster

Particularly sensitive:

Stars that start burning helium after becoming red giants

$$T_{\text{HB}} \sim 10 \text{ keV}$$

Axion emission reduces lifetime of
HB stars

\Rightarrow Observations require
 $g_{ax} \lesssim 10^{-10} \text{ GeV}$ for $m_a \lesssim 10 \text{ keV}$

\Rightarrow Excludes QCD axions with
 $m_a \gtrsim 10 \text{ eV}$

5.5.4 Cooling hints

White dwarfs have electron-degenerate
cores

\Rightarrow efficient axion emission if
 $g_{ae} \neq 0$

Can measure energy loss via change of
pulsation period (\dot{P}/P)

\Rightarrow Several systems are found
to cool more quickly
than expected

\hookrightarrow Hint for axions?

Requires $f_a \sim 10^9 \text{ GeV}$
 $m_a \sim 10 \text{ meV}$

\Rightarrow outside of
preferred DM
window

5.5.5 Supernova 1987A

Mechanism: Core-collapse SN

↳ Nuclear fusion insufficient to provide pressure against gravity

↳ Core heats up to MeV temperatures

⇒ Significant fraction of rest mass converted into neutrino pulse

⇒ Visible light is only emitted once shock wave reaches stellar surface

SN 1987a: Distance: 50 kpc

→ 25 neutrino events observed within ~ 10 seconds across 3 experiments

→ Visible light only observed several hours later

If axions exist, they can be produced in the SN core and be converted into photons in the Galactic magnetic field

\Rightarrow γ -ray burst coincident with
neutrinos

Conversion again requires coherence:

$$q \cdot \ell < 1$$

$$E_a \sim 50 \text{ MeV}$$

$$\ell \sim 50 \text{ kpc}$$

$$\Rightarrow m_a \lesssim 10^{-10} \text{ eV}$$

Very strong bound for the ultralight
axion-like particles:

$$g_{\text{ax}} \lesssim 5 \cdot 10^{-12} \text{ GeV}^{-1}$$

Note: Even if axions are
not converted, one finds

$$g_{\text{ax}} \lesssim 6 \cdot 10^{-9} \text{ GeV}^{-1}$$

from heat loss

5.6 Fuzzy dark matter

In lecture 1: Determined shapes of DM halos by solving eq. of hydrostatic equilibrium

$$\frac{k_B T}{G m} \frac{d}{dr} \left(\frac{r^2}{\rho} \frac{d\rho}{dr} \right) = 4\pi r^2 \rho(r)$$

What if quantum effects become important?

$$\lambda_{dB} = \frac{h}{m_{DM} \sqrt{\frac{R}{M(R) G}}}$$

$$\lambda_{dB} \sim \text{kpc} \quad \text{for} \quad m_{DM} \sim 10^{-22} \text{ eV}$$

For QCD axions: $m_a = \frac{m_\pi f_\pi}{2 f_a}$

$$\geq \frac{m_\pi f_\pi}{2 M_{\text{pe}}} \sim 5 \cdot 10^{-13} \text{ eV}$$

But axion-like particles can have much smaller masses.

\Rightarrow Need to replace $\rho \rightarrow m |\psi|^2$
where ψ is the axion wave function

\Rightarrow Gravitational potential Φ given by

$$\nabla^2 \Phi = 4\pi G m |\psi|^2$$

(103)

ψ must satisfy Schrödinger equation

$$-\frac{\hbar^2}{2m} \nabla^2 \psi + m \bar{\phi} \psi = m E \psi$$

Numerical solutions imply upper bound on central densities of DM halos:

$$\rho_c \leq 300 \frac{\text{GeV}}{\text{cm}^3} \left(\frac{m_{\text{DM}}}{10^{-22} \text{eV}} \right)^6 \left(\frac{M_{\text{halo}}}{10^5 M_\odot} \right)^4$$

and lower bounds on their size

$$r \gtrsim 0.3 \text{ kpc} \left(\frac{m_{\text{DM}}}{10^{-22} \text{eV}} \right)^{-2} \left(\frac{M_{\text{halo}}}{10^5 M_\odot} \right)^{-1}$$

\Rightarrow Axion-like particles with $m_a \leq 10^{-22} \text{eV}$ cannot be all of DM

But for $m_a \sim 10^{-22} - 10^{-21} \text{eV}$ agreement better than for CDM!