

The perturbative gradient flow

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What is the gradient flow?
Perturbative solution
Effective Field Theories
Calculational techniques

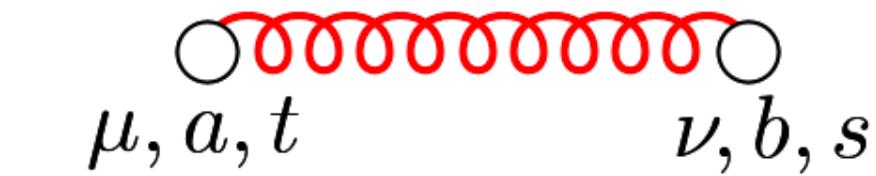
Part 2

Feynman rules

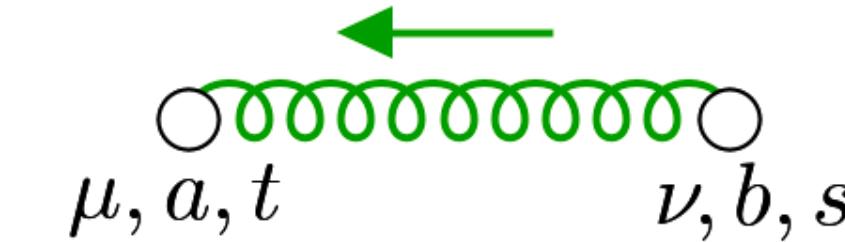
$$\mathcal{L} = \mathcal{L}_{\text{QCD}} + \mathcal{L}_B + \mathcal{L}_\chi$$

$$\mathcal{L}_B \sim \int_0^\infty dt \, L_\mu \left(\partial_t B_\mu - \mathcal{D}_\nu G_{\nu\mu} \right)$$

L_μ Lagrange multiplier field



$$\frac{\delta^{ab}}{p^2} \left(\delta_{\mu\nu} - \xi \frac{p_\mu p_\nu}{p^2} \right) e^{-(t+s)p^2}$$



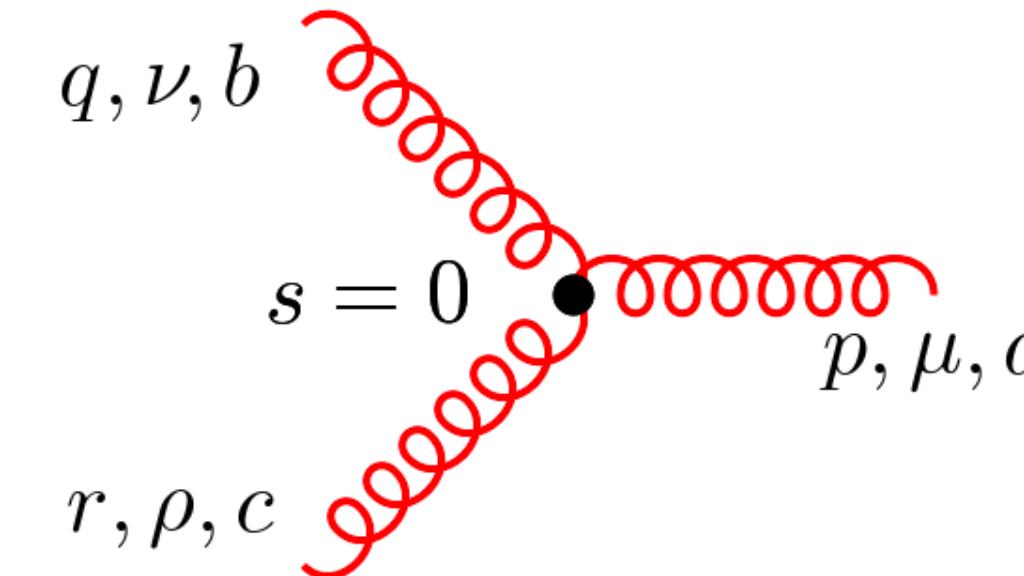
$$\delta_{ab} \delta_{\mu\nu} \theta(t-s) e^{-(t-s)p^2}$$

“gluon flow line”

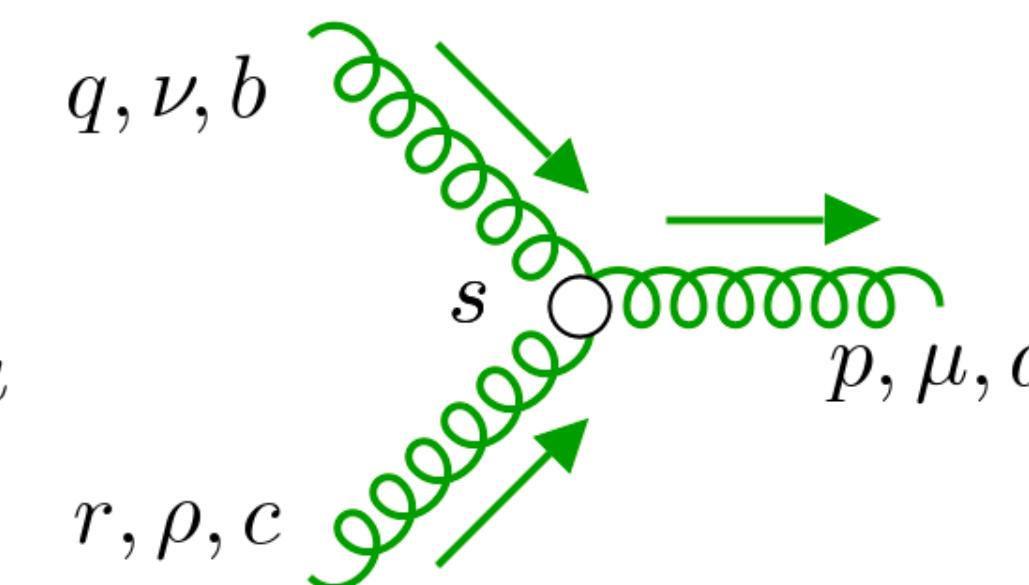
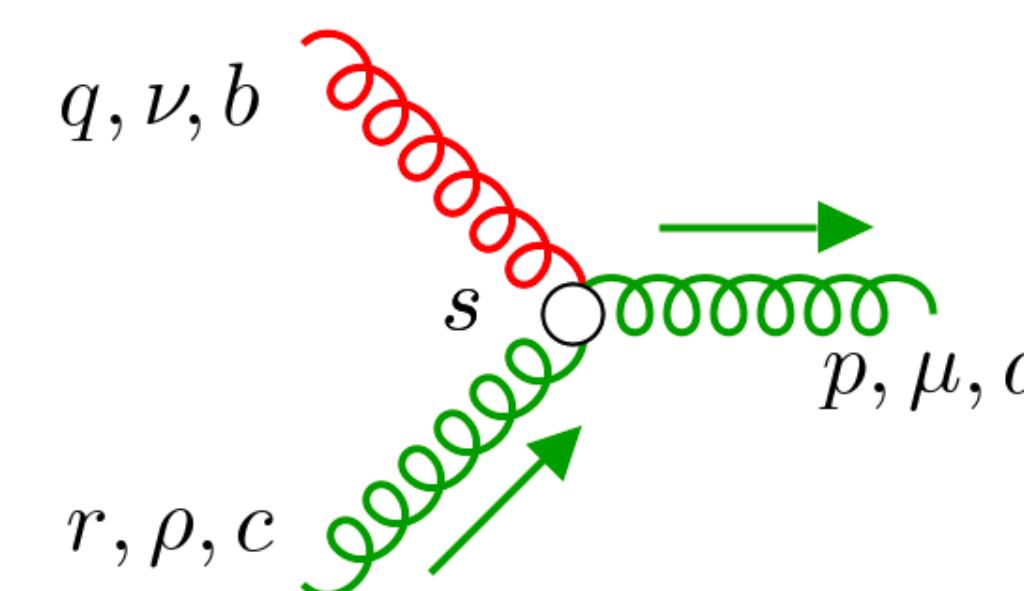
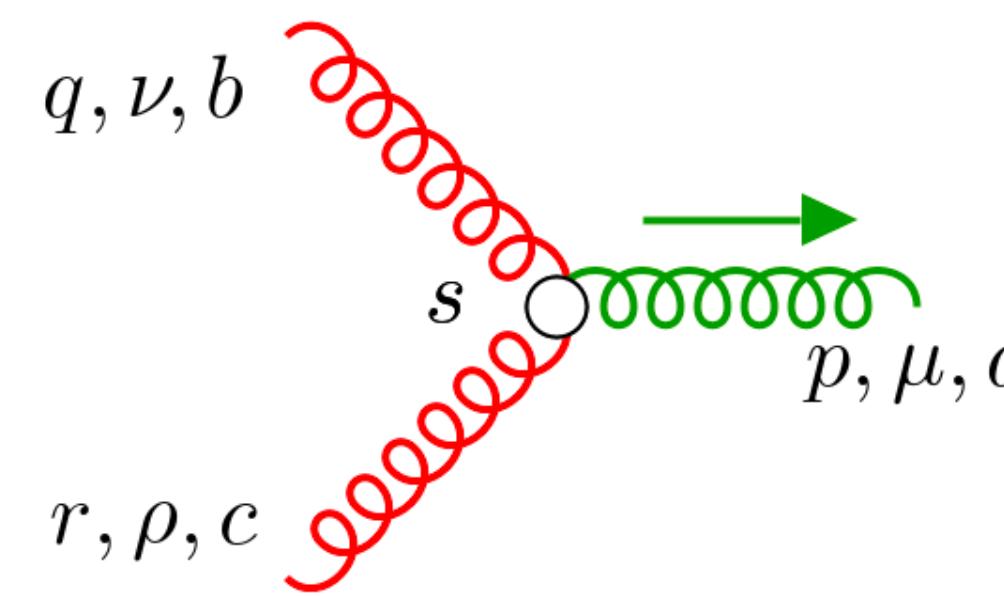
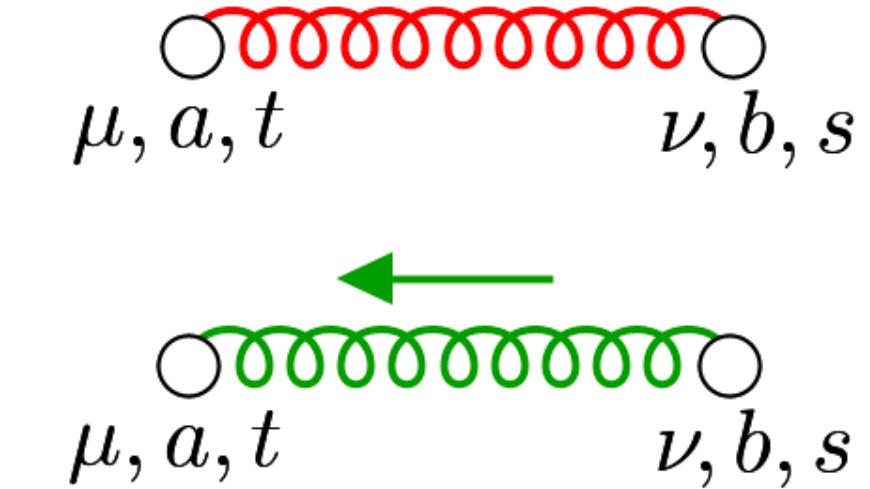
analogously for quarks:

$$\mathcal{L}_\chi \sim \int_0^\infty dt \, \bar{\lambda} (\partial_t - \Delta) \lambda + \text{h.c.}$$

Vertices



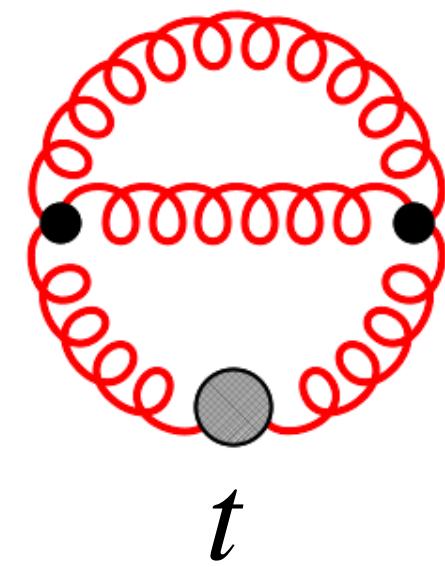
regular 3-gluon vertex



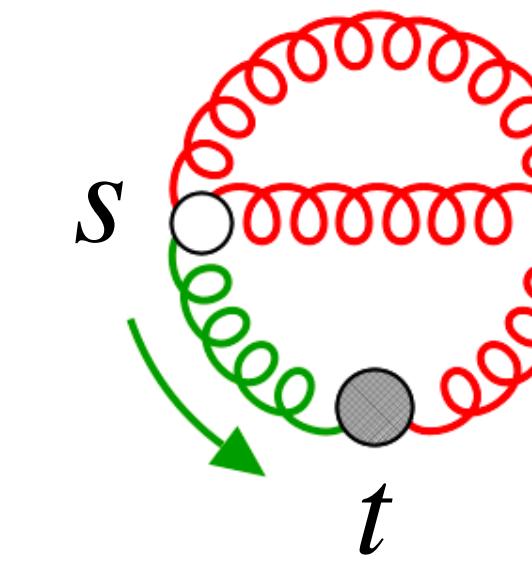
$$-igf^{abc} \int_0^\infty ds (\delta_{\nu\rho}(r-q)_\mu + 2\delta_{\mu\nu}q_\rho - 2\delta_{\mu\rho}r_\nu + (\kappa-1)(\delta_{\mu\rho}q_\nu - \delta_{\mu\nu}r_\rho))$$

analogously for 4-gluon vertex and quarks

Higher orders



$$\sim \int_p \int_k \frac{e^{-2\textcolor{red}{t} p^2}}{p^4 k^2 (p - k)^2}$$



$$\int_0^t \textcolor{red}{ds} \int_p \int_k \frac{e^{-(2t-s)p^2}}{p^2 k^2 (p - k)^2}$$

- more loop integrals
- integration over flow-time parameters
- renormalization: same as fundamental QCD!

Integration-by-parts relations

- After tensor reduction, we end up with many scalar integrals of the form

$$I(\{t_f^{\text{up}}\}, \{T_i\}, \{a_i\}) = \left(\prod_{f=1}^F \int_0^{t_f^{\text{up}}} dt_f \right) \int_{k_1, \dots, k_L} \frac{\exp[-(T_1 q_1^2 + \dots + T_N q_N^2)]}{q_1^{2a_1} \cdots q_N^{2a_N}}$$

with q_i linear combinations of k_j and T_i linear combinations of t_j , e.g. $q_1 = k_1 - k_2$ and $T_1 = t + 2t_1 - t_3$

- Chetyrkin and Tkachov observed [Tkachov 1981; Chetyrkin, Tkachov 1981]

$$\int_{k_1, \dots, k_L} \frac{\partial}{\partial k_i^\mu} \left(\tilde{q}_j^\mu \frac{1}{P_1^{a_1} \cdots P_N^{a_N}} \right) = 0$$

- ⇒ Linear relations between Feynman integrals
- Can easily be adopted to gradient-flow integrals
 - Additional new relations for gradient-flow integrals: [Artz, RH, Lange, Neumann, Prausa '19]

$$\int_0^{t_f^{\text{up}}} dt_f \partial_{t_f} F(t_f, \dots) = F(t_f^{\text{up}}, \dots) - F(0, \dots)$$

Laporta algorithm

- Schematically integration-by-parts read

$$0 = (d - a_1) I(a_1, a_2, a_3) + (a_1 - a_2) I(a_1 + 1, a_2 - 1, a_3) + (2a_3 + a_1 - a_2) I(a_1 + 1, a_2, a_3 - 1)$$

- Rarely possible to find general solution like

$$I(a_1, a_2, a_3) = a_1 I(a_1 - 1, a_2, a_3) + (d + a_1 - a_2) I(a_1, a_2 - 1, a_3) + 2a_3 I(a_1, a_2, a_3 - 1)$$

- Instead set up system of equations and solve it [Laporta 2000] :

- Insert seeds $\{a_1 = 1, a_2 = 1, a_3 = 1\}$, $\{a_1 = 2, a_2 = 1, a_3 = 1\}$, ...:

$$0 = (d - 1) I(1, 1, 1) + I(2, 1, 0),$$

$$0 = (d - 2) I(2, 1, 1) + I(3, 0, 1) - I(3, 1, 0),$$

⋮

- Solve with Gaussian elimination

⇒ Express integrals through significantly smaller number of master integrals

Laporta algorithm

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$$0 = (d - a_1) I(a_1, a_2, a_3) + (a_1 - a_2) I(a_1 + 1, a_2 - 1, a_3) + (2a_3 + a_1 - a_2) I(a_1 + 1, a_2, a_3 - 1)$$

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- Instead set up system of equations and solve it [Laporta 2000] :

- Insert seeds $\{a_1 = 1, a_2 = 1, a_3 = 1\}$, $\{a_1 = 2, a_2 = 1, a_3 = 1\}$, ...:

$$0 = (d - 1) I(1, 1, 1) + I(2, 1, 0),$$

$$0 = (d - 2) I(2, 1, 1) + I(3, 0, 1) - I(3, 1, 0),$$

e.g. NNLO chromo-magnetic dipole operator:
 O(4000) integrals reduced to 13 master integrals

- Solve with Gaussian elimination

⇒ Express integrals through significantly smaller number of master integrals

Numerical evaluation of the master integrals

$$\int_0^1 du_1 u_1^{c_1} \dots \int_0^1 du_f u_f^{c_f} \iint_{p_1, p_2, p_3} \frac{\exp\left(-\mathbf{p}^T A(u_1, \dots, u_f) \mathbf{p}\right)}{p_1^2 p_2^2 p_3^2 (p_1 - p_2)^2 (p_1 - p_3)^2 (p_2 - p_3)^2}$$

Numerical evaluation of the master integrals

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Schwinger parameters:

$$\frac{1}{p^2} = \int_0^\infty dx e^{-x p^2}$$

Numerical evaluation of the master integrals

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Schwinger parameters:

$$\frac{1}{p^2} = \int_0^\infty dx e^{-x p^2}$$

$$\int_0^1 du_1 u_1^{c_1} \dots \int_0^1 du_f u_f^{c_f} \int_0^\infty dx_1 \dots \int_0^\infty dx_6 \iint_{p_1, p_2, p_3} \exp(-\mathbf{p}^T B(u_1, \dots, u_f, x_1, \dots, x_6) \mathbf{p})$$

Numerical evaluation of the master integrals

$$\int_0^1 du_1 u_1^{c_1} \cdots \int_0^1 du_f u_f^{c_f} \iint_{p_1, p_2, p_3} \frac{\exp(-\mathbf{p}^T A(u_1, \dots, u_f) \mathbf{p})}{p_1^2 p_2^2 p_3^2 (p_1 - p_2)^2 (p_1 - p_3)^2 (p_2 - p_3)^2}$$

Schwinger parameters:

$$\frac{1}{p^2} = \int_0^\infty dx e^{-x p^2}$$

$$\int_0^1 du_1 u_1^{c_1} \cdots \int_0^1 du_f u_f^{c_f} \int_0^\infty dx_1 \cdots \int_0^\infty dx_6 \iint_{p_1, p_2, p_3} \exp(-\mathbf{p}^T B(u_1, \dots, u_f, x_1, \dots, x_6) \mathbf{p})$$

$$\int_0^1 du_1 u_1^{c_1} \cdots \int_0^1 du_f u_f^{c_f} \int_0^\infty dx_1 \cdots \int_0^\infty dx_6 [\det B(u_1, \dots, u_f, x_1, \dots, x_6)]^{-D/2}$$

$$\int_0^1 du_1 u_1 \int_0^1 du_2 \int_{p_1, p_2, p_3} \frac{\exp(-p_1^2 - u_1 p_2^2 - u_1 u_2 p_3^2 - 2(p_1 - p_2)^2)}{p_1^2 p_2^2 p_3^2 (p_1 - p_2)^2 (p_1 - p_3)^2 (p_2 - p_3)^2} \longrightarrow$$

$$\int_0^1 du_1 u_1 \int_0^1 du_2 \int_{p_1, p_2, p_3} \frac{\exp(-p_1^2 - u_1 p_2^2 - u_1 u_2 p_3^2 - 2(p_1 - p_2)^2)}{p_1^2 p_2^2 p_3^2 (p_1 - p_2)^2 (p_1 - p_3)^2 (p_2 - p_3)^2} \longrightarrow$$

$$\int_0^1 du_1 \int_0^1 du_2 \int_0^\infty dx_1 \int_0^\infty dx_2 \int_0^\infty dx_3 \int_0^\infty dx_4 \int_0^\infty dx_5 \int_0^\infty dx_6$$

$$\int_0^1 du_1 u_1 \int_0^1 du_2 \int_{p_1, p_2, p_3} \frac{\exp(-p_1^2 - u_1 p_2^2 - u_1 u_2 p_3^2 - 2(p_1 - p_2)^2)}{p_1^2 p_2^2 p_3^2 (p_1 - p_2)^2 (p_1 - p_3)^2 (p_2 - p_3)^2} \longrightarrow$$

$$\begin{aligned} & \int_0^1 du_1 \int_0^1 du_2 \int_0^\infty dx_1 \int_0^\infty dx_2 \int_0^\infty dx_3 \int_0^\infty dx_4 \int_0^\infty dx_5 \int_0^\infty dx_6 \quad u_1 x_1^{-\epsilon} x_2^{-\epsilon} x_3^{-\epsilon} x_4^{-\epsilon} x_5^{-\epsilon} x_6^{-\epsilon} (3 u_1^2 u_2 x_1 x_2 x_3 x_4 x_5 x_6 + u_1^2 u_2 x_1 x_2 x_3 x_4 x_6 + \\ & + u_1^2 u_2 x_1 x_2 x_3 x_5 x_6 + u_1^2 u_2 x_2 x_3 x_4 x_5 x_6 + 2 u_1 u_2 x_1 x_2 x_3 x_4 x_5 x_6 + \\ & + 3 u_1 u_2 x_1 x_2 x_3 x_4 x_5 + 2 u_1 u_2 x_1 x_2 x_3 x_4 x_6 + u_1 u_2 x_1 x_2 x_3 x_4 + \\ & + u_1 u_2 x_1 x_2 x_3 x_5 x_6 + u_1 u_2 x_1 x_2 x_3 x_5 + u_1 u_2 x_1 x_2 x_3 x_6 + \\ & + 3 u_1 u_2 x_1 x_3 x_4 x_5 x_6 + u_1 u_2 x_1 x_3 x_4 x_6 + u_1 u_2 x_1 x_3 x_5 x_6 + \\ & + 2 u_1 u_2 x_2 x_3 x_4 x_5 x_6 + u_1 u_2 x_2 x_3 x_4 x_5 + u_1 u_2 x_2 x_3 x_5 x_6 + \\ & + u_1 u_2 x_3 x_4 x_5 x_6 + 3 u_1 x_1 x_2 x_3 x_4 x_5 + 3 u_1 x_1 x_2 x_3 x_4 x_6 + \\ & + u_1 x_1 x_2 x_3 x_4 + u_1 x_1 x_2 x_3 x_5 + u_1 x_1 x_2 x_3 x_6 + 3 u_1 x_1 x_2 x_4 x_5 x_6 + \\ & + u_1 x_1 x_2 x_4 x_6 + u_1 x_1 x_2 x_5 x_6 + u_1 x_2 x_3 x_4 x_5 + u_1 x_2 x_3 x_4 x_6 + \\ & + u_1 x_2 x_4 x_5 x_6 + 2 x_1 x_2 x_3 x_4 x_5 + 2 x_1 x_2 x_3 x_4 x_6 + x_1 x_2 x_3 x_4 + \\ & + x_1 x_2 x_3 x_5 + x_1 x_2 x_3 x_6 + 2 x_1 x_2 x_4 x_5 x_6 + 3 x_1 x_2 x_4 x_5 + \\ & + 2 x_1 x_2 x_4 x_6 + x_1 x_2 x_4 + x_1 x_2 x_5 x_6 + x_1 x_2 x_5 + x_1 x_2 x_6 + \\ & + 3 x_1 x_3 x_4 x_5 + 3 x_1 x_3 x_4 x_6 + x_1 x_3 x_4 + x_1 x_3 x_5 + x_1 x_3 x_6 + \\ & + 3 x_1 x_4 x_5 x_6 + x_1 x_4 x_6 + x_1 x_5 x_6 + 2 x_2 x_3 x_4 x_5 + 2 x_2 x_3 x_4 x_6 + \\ & + x_2 x_3 x_4 + x_2 x_3 x_5 + x_2 x_3 x_6 + 2 x_2 x_4 x_5 x_6 + \\ & + x_2 x_4 x_5 + x_2 x_5 x_6 + x_3 x_4 x_5 + x_3 x_4 x_6 + x_4 x_5 x_6)^{\epsilon-2} \end{aligned}$$

$$\int_0^1 du_1 u_1 \int_0^1 du_2 \int_{p_1, p_2, p_3} \frac{\exp(-p_1^2 - u_1 p_2^2 - u_1 u_2 p_3^2 - 2(p_1 - p_2)^2)}{p_1^2 p_2^2 p_3^2 (p_1 - p_2)^2 (p_1 - p_3)^2 (p_2 - p_3)^2} \longrightarrow$$

$$\int_0^1 du_1 \int_0^1 du_2 \int_0^\infty dx_1 \int_0^\infty dx_2 \int_0^\infty dx_3 \int_0^\infty dx_4 \int_0^\infty dx_5 \int_0^\infty dx_6 \quad u_1 x_1^{-\epsilon} x_2^{-\epsilon} x_3^{-\epsilon} x_4^{-\epsilon} x_5^{-\epsilon} x_6^{-\epsilon} (3 u_1^2 u_2 x_1 x_2 x_3 x_4 x_5 x_6 + u_1^2 u_2 x_1 x_2 x_3 x_4 x_6 + \\ + u_1^2 u_2 x_1 x_2 x_3 x_5 x_6 + u_1^2 u_2 x_2 x_3 x_4 x_5 x_6 + 2 u_1 u_2 x_1 x_2 x_3 x_4 x_5 x_6 + \\ + 3 u_1 u_2 x_1 x_2 x_3 x_4 x_5 + 2 u_1 u_2 x_1 x_2 x_3 x_4 x_6 + u_1 u_2 x_1 x_2 x_3 x_4 + \\ + u_1 u_2 x_1 x_2 x_3 x_5 x_6 + u_1 u_2 x_1 x_2 x_3 x_5 + u_1 u_2 x_1 x_2 x_3 x_6 + \\ + 3 u_1 u_2 x_1 x_3 x_4 x_5 x_6 + u_1 u_2 x_1 x_3 x_4 x_6 + u_1 u_2 x_1 x_3 x_5 x_6 + \\ + 2 u_1 u_2 x_2 x_3 x_4 x_5 x_6 + u_1 u_2 x_2 x_3 x_4 x_5 + u_1 u_2 x_2 x_3 x_5 x_6 + \\ + u_1 u_2 x_3 x_4 x_5 x_6 + 3 u_1 x_1 x_2 x_3 x_4 x_5 + 3 u_1 x_1 x_2 x_3 x_4 x_6 + \\ + u_1 x_1 x_2 x_3 x_4 + u_1 x_1 x_2 x_3 x_5 + u_1 x_1 x_2 x_3 x_6 + 3 u_1 x_1 x_2 x_4 x_5 x_6 + \\ + u_1 x_1 x_2 x_4 x_6 + u_1 x_1 x_2 x_5 x_6 + u_1 x_2 x_3 x_4 x_5 + u_1 x_2 x_3 x_4 x_6 + \\ + u_1 x_2 x_4 x_5 x_6 + 2 x_1 x_2 x_3 x_4 x_5 + 2 x_1 x_2 x_3 x_4 x_6 + x_1 x_2 x_3 x_4 + \\ + x_1 x_2 x_3 x_5 + x_1 x_2 x_3 x_6 + 2 x_1 x_2 x_4 x_5 x_6 + 3 x_1 x_2 x_4 x_5 + \\ + 2 x_1 x_2 x_4 x_6 + x_1 x_2 x_4 + x_1 x_2 x_5 x_6 + x_1 x_2 x_5 + x_1 x_2 x_6 + \\ + 3 x_1 x_3 x_4 x_5 + 3 x_1 x_3 x_4 x_6 + x_1 x_3 x_4 + x_1 x_3 x_5 + x_1 x_3 x_6 + \\ + 3 x_1 x_4 x_5 x_6 + x_1 x_4 x_6 + x_1 x_5 x_6 + 2 x_2 x_3 x_4 x_5 + 2 x_2 x_3 x_4 x_6 + \\ + x_2 x_3 x_4 + x_2 x_3 x_5 + x_2 x_3 x_6 + 2 x_2 x_4 x_5 x_6 + \\ + x_2 x_4 x_5 + x_2 x_5 x_6 + x_3 x_4 x_5 + x_3 x_4 x_6 + x_4 x_5 x_6)^{\epsilon-2}$$
$$\int_0^\infty dx f(x) = \int_0^1 \frac{dy}{y^2} f(x(y))$$

$$\int_0^1 du_1 u_1 \int_0^1 du_2 \int_{p_1, p_2, p_3} \frac{\exp(-p_1^2 - u_1 p_2^2 - u_1 u_2 p_3^2 - 2(p_1 - p_2)^2)}{p_1^2 p_2^2 p_3^2 (p_1 - p_2)^2 (p_1 - p_3)^2 (p_2 - p_3)^2} \longrightarrow$$

$$\int_0^1 du_1 \int_0^1 du_2 \int_0^\infty dx_1 \int_0^\infty dx_2 \int_0^\infty dx_3 \int_0^\infty dx_4 \int_0^\infty dx_5 \int_0^\infty dx_6 \quad u_1 x_1^{-\epsilon} x_2^{-\epsilon} x_3^{-\epsilon} x_4^{-\epsilon} x_5^{-\epsilon} x_6^{-\epsilon} (3 u_1^2 u_2 x_1 x_2 x_3 x_4 x_5 x_6 + u_1^2 u_2 x_1 x_2 x_3 x_4 x_6 + \\ + u_1^2 u_2 x_1 x_2 x_3 x_5 x_6 + u_1^2 u_2 x_2 x_3 x_4 x_5 x_6 + 2 u_1 u_2 x_1 x_2 x_3 x_4 x_5 x_6 + \\ + 3 u_1 u_2 x_1 x_2 x_3 x_4 x_5 + 2 u_1 u_2 x_1 x_2 x_3 x_4 x_6 + u_1 u_2 x_1 x_2 x_3 x_4 + \\ + u_1 u_2 x_1 x_2 x_3 x_5 x_6 + u_1 u_2 x_1 x_2 x_3 x_5 + u_1 u_2 x_1 x_2 x_3 x_6 + \\ + 3 u_1 u_2 x_1 x_3 x_4 x_5 x_6 + u_1 u_2 x_1 x_3 x_4 x_6 + u_1 u_2 x_1 x_3 x_5 x_6 + \\ + 2 u_1 u_2 x_2 x_3 x_4 x_5 x_6 + u_1 u_2 x_2 x_3 x_4 x_5 + u_1 u_2 x_2 x_3 x_5 x_6 + \\ + u_1 u_2 x_3 x_4 x_5 x_6 + 3 u_1 x_1 x_2 x_3 x_4 x_5 + 3 u_1 x_1 x_2 x_3 x_4 x_6 + \\ + u_1 x_1 x_2 x_3 x_4 + u_1 x_1 x_2 x_3 x_5 + u_1 x_1 x_2 x_3 x_6 + 3 u_1 x_1 x_2 x_4 x_5 x_6 + \\ + u_1 x_1 x_2 x_4 x_6 + u_1 x_1 x_2 x_5 x_6 + u_1 x_2 x_3 x_4 x_5 + u_1 x_2 x_3 x_4 x_6 + \\ + u_1 x_2 x_4 x_5 x_6 + 2 x_1 x_2 x_3 x_4 x_5 + 2 x_1 x_2 x_3 x_4 x_6 + x_1 x_2 x_3 x_4 + \\ + x_1 x_2 x_3 x_5 + x_1 x_2 x_3 x_6 + 2 x_1 x_2 x_4 x_5 x_6 + 3 x_1 x_2 x_4 x_5 + \\ + 2 x_1 x_2 x_4 x_6 + x_1 x_2 x_4 + x_1 x_2 x_5 x_6 + x_1 x_2 x_5 + x_1 x_2 x_6 + \\ + 3 x_1 x_3 x_4 x_5 + 3 x_1 x_3 x_4 x_6 + x_1 x_3 x_4 + x_1 x_3 x_5 + x_1 x_3 x_6 + \\ + 3 x_1 x_4 x_5 x_6 + x_1 x_4 x_6 + x_1 x_5 x_6 + 2 x_2 x_3 x_4 x_5 + 2 x_2 x_3 x_4 x_6 + \\ + x_2 x_3 x_4 + x_2 x_3 x_5 + x_2 x_3 x_6 + 2 x_2 x_4 x_5 x_6 + \\ + x_2 x_4 x_5 + x_2 x_5 x_6 + x_3 x_4 x_5 + x_3 x_4 x_6 + x_4 x_5 x_6)^{\epsilon-2}$$

$y = \frac{1}{1+x}$

$$\int_0^\infty dx f(x) = \int_0^1 \frac{dy}{y^2} f(x(y))$$

overlapping singularities
as $x_i, u_j \rightarrow 0$

$$\int_0^1 du_1 u_1 \int_0^1 du_2 \int_{p_1, p_2, p_3} \frac{\exp(-p_1^2 - u_1 p_2^2 - u_1 u_2 p_3^2 - 2(p_1 - p_2)^2)}{p_1^2 p_2^2 p_3^2 (p_1 - p_2)^2 (p_1 - p_3)^2 (p_2 - p_3)^2} \longrightarrow$$

$$\int_0^1 du_1 \int_0^1 du_2 \int_0^\infty dx_1 \int_0^\infty dx_2 \int_0^\infty dx_3 \int_0^\infty dx_4 \int_0^\infty dx_5 \int_0^\infty dx_6 \quad u_1 x_1^{-\epsilon} x_2^{-\epsilon} x_3^{-\epsilon} x_4^{-\epsilon} x_5^{-\epsilon} x_6^{-\epsilon} (3 u_1^2 u_2 x_1 x_2 x_3 x_4 x_5 x_6 + u_1^2 u_2 x_1 x_2 x_3 x_4 x_6 + \\ + u_1^2 u_2 x_1 x_2 x_3 x_5 x_6 + u_1^2 u_2 x_2 x_3 x_4 x_5 x_6 + 2 u_1 u_2 x_1 x_2 x_3 x_4 x_5 x_6 + \\ + 3 u_1 u_2 x_1 x_2 x_3 x_4 x_5 + 2 u_1 u_2 x_1 x_2 x_3 x_4 x_6 + u_1 u_2 x_1 x_2 x_3 x_4 + \\ + u_1 u_2 x_1 x_2 x_3 x_5 x_6 + u_1 u_2 x_1 x_2 x_3 x_5 + u_1 u_2 x_1 x_2 x_3 x_6 + \\ + 3 u_1 u_2 x_1 x_3 x_4 x_5 x_6 + u_1 u_2 x_1 x_3 x_4 x_6 + u_1 u_2 x_1 x_3 x_5 x_6 + \\ + 2 u_1 u_2 x_2 x_3 x_4 x_5 x_6 + u_1 u_2 x_2 x_3 x_4 x_5 + u_1 u_2 x_2 x_3 x_5 x_6 + \\ + u_1 u_2 x_3 x_4 x_5 x_6 + 3 u_1 x_1 x_2 x_3 x_4 x_5 + 3 u_1 x_1 x_2 x_3 x_4 x_6 + \\ + u_1 x_1 x_2 x_3 x_4 + u_1 x_1 x_2 x_3 x_5 + u_1 x_1 x_2 x_3 x_6 + 3 u_1 x_1 x_2 x_4 x_5 x_6 + \\ + u_1 x_1 x_2 x_4 x_6 + u_1 x_1 x_2 x_5 x_6 + u_1 x_2 x_3 x_4 x_5 + u_1 x_2 x_3 x_4 x_6 + \\ + u_1 x_2 x_4 x_5 x_6 + 2 x_1 x_2 x_3 x_4 x_5 + 2 x_1 x_2 x_3 x_4 x_6 + x_1 x_2 x_3 x_4 + \\ + x_1 x_2 x_3 x_5 + x_1 x_2 x_3 x_6 + 2 x_1 x_2 x_4 x_5 x_6 + 3 x_1 x_2 x_4 x_5 + \\ + 2 x_1 x_2 x_4 x_6 + x_1 x_2 x_4 + x_1 x_2 x_5 x_6 + x_1 x_2 x_5 + x_1 x_2 x_6 + \\ + 3 x_1 x_3 x_4 x_5 + 3 x_1 x_3 x_4 x_6 + x_1 x_3 x_4 + x_1 x_3 x_5 + x_1 x_3 x_6 + \\ + 3 x_1 x_4 x_5 x_6 + x_1 x_4 x_6 + x_1 x_5 x_6 + 2 x_2 x_3 x_4 x_5 + 2 x_2 x_3 x_4 x_6 + \\ + x_2 x_3 x_4 + x_2 x_3 x_5 + x_2 x_3 x_6 + 2 x_2 x_4 x_5 x_6 + \\ + x_2 x_4 x_5 + x_2 x_5 x_6 + x_3 x_4 x_5 + x_3 x_4 x_6 + x_4 x_5 x_6)^{\epsilon-2}$$

$y = \frac{1}{1+x}$

$$\int_0^\infty dx f(x) = \int_0^1 \frac{dy}{y^2} f(x(y))$$

overlapping singularities
as $x_i, u_j \rightarrow 0$

→ sector decomposition

example:
[Heinrich '08]

$$I = \int_0^1 dx \int_0^1 dy x^{-1-a\epsilon} y^{-b\epsilon} \left(x + (1-x)y \right)^{-1}$$

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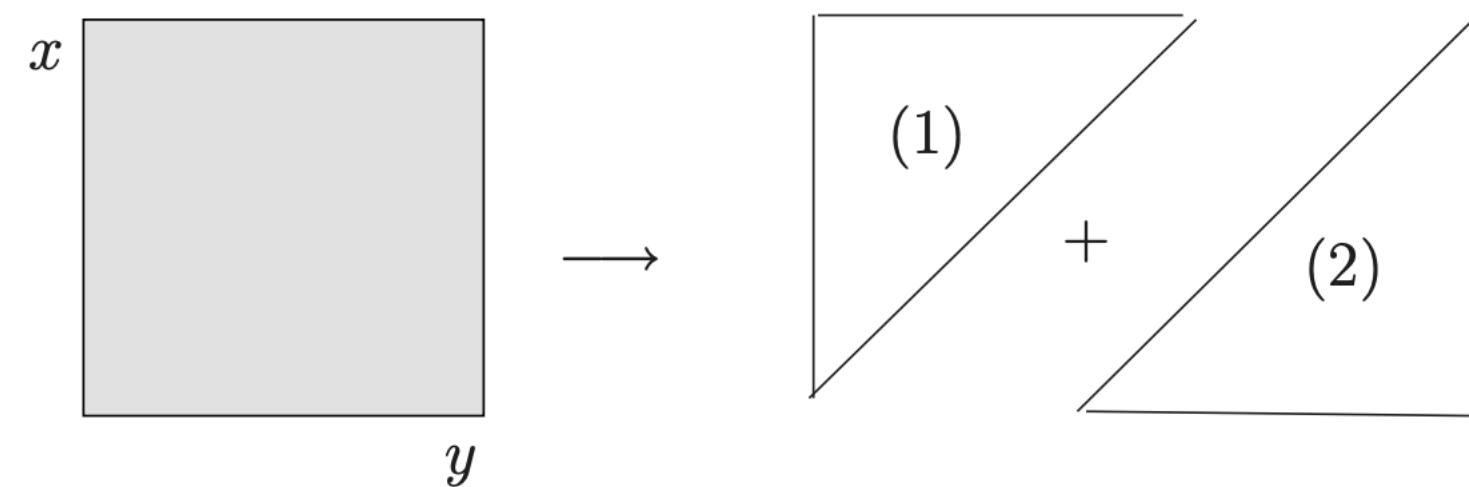
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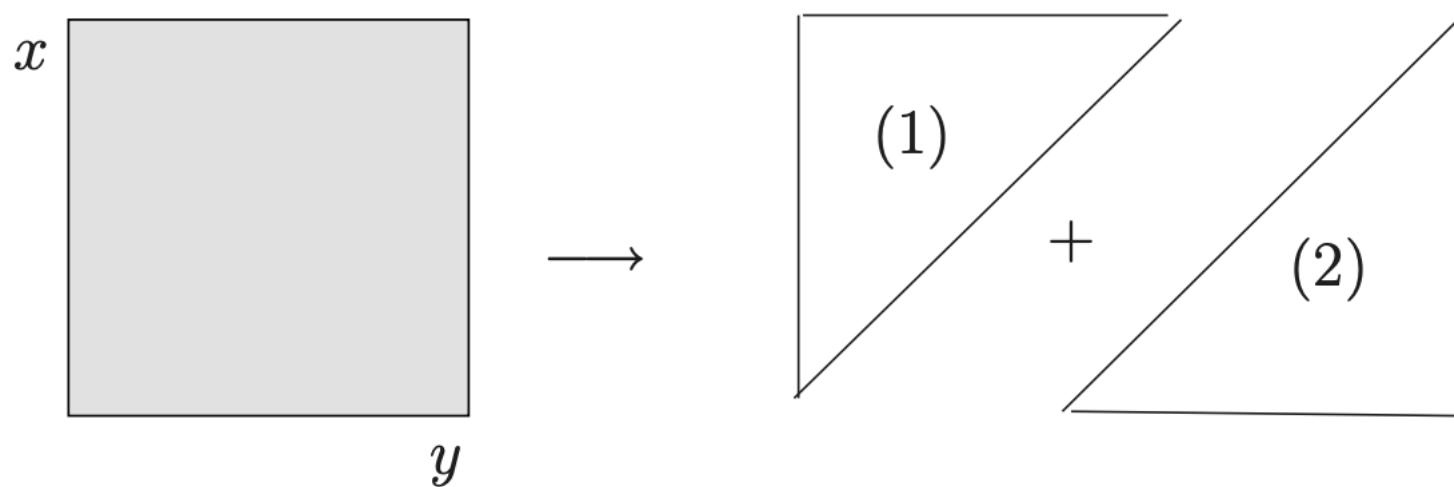
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&= -\frac{1}{(a+b)\epsilon} \int_0^1 dt \frac{t^{-b\epsilon}}{1+t}
\end{aligned}$$

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\end{aligned}$$

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\end{aligned}$$

$$\int_0^1 dt \left(\frac{\ln^n t}{t}\right)_+ f(t) = \int_0^1 dt \frac{\ln^n t}{t} [f(t) - f(0)]$$

pySecDec [Borowka, Heinrich, Jahn, Jones, Kerner, Schlenk, Zirke '18]

$$\int_0^1 du_1 u_1 \int_0^1 du_2 \int_{p_1, p_2, p_3} \frac{\exp(-p_1^2 - u_1 p_2^2 - u_1 u_2 p_3^2 - 2(p_1 - p_2)^2)}{p_1^2 p_2^2 p_3^2 (p_1 - p_2)^2 (p_1 - p_3)^2 (p_2 - p_3)^2} = \frac{1}{(4\pi)^{3d/2}} \left[+ \text{ep}^(-1) * ((-1.20205690407937649) + (6.74709950249940753e-9) * \text{numerr}) + \text{ep}^0 * ((-11.4409624237256917) + (4.99888756503079786e-8) * \text{numerr}) \right]$$

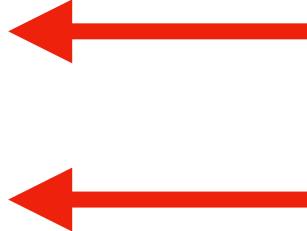
[RH, Nellopolous '22 (unpublished)]
earlier work: [RH, Neumann '16]

Effective Field Theories

$$\mathcal{L} = \mathcal{L}^{\leq 4} + \sum_{d>4} \frac{1}{\Lambda^{d-4}} \sum_i C_i^{(d)} \mathcal{O}_i^{(d)}$$

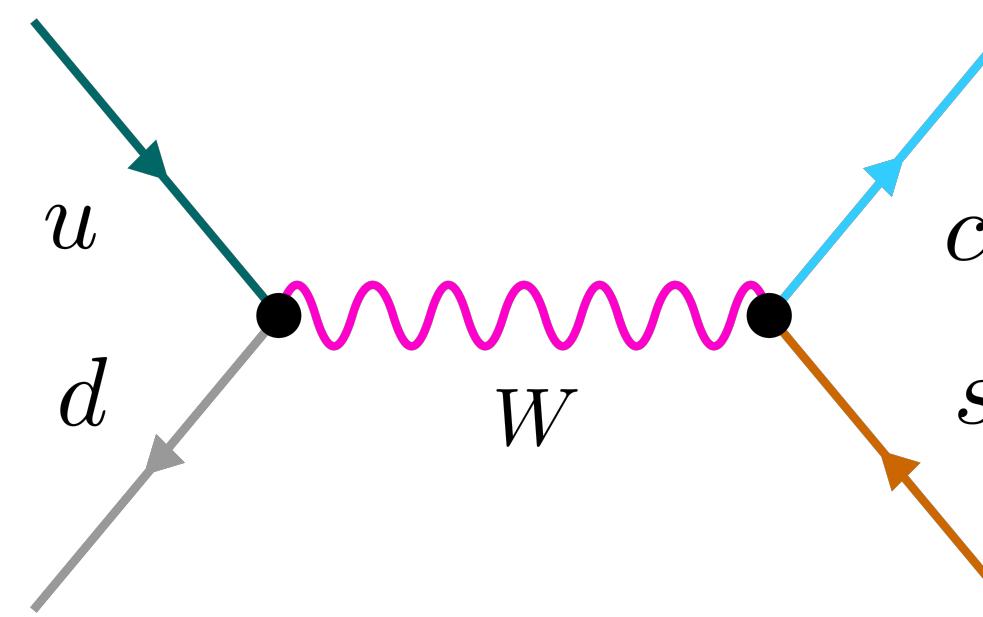
some problems:

- many operators (SMEFT: 2499 @ dim 6)
- get a non-redundant basis (EoMs, IbP, Fierz, Schouten, ...)
- determine $C_i^{(d)}$
- determine $\langle \mathcal{O}_i^{(d)} \rangle$
- renormalization

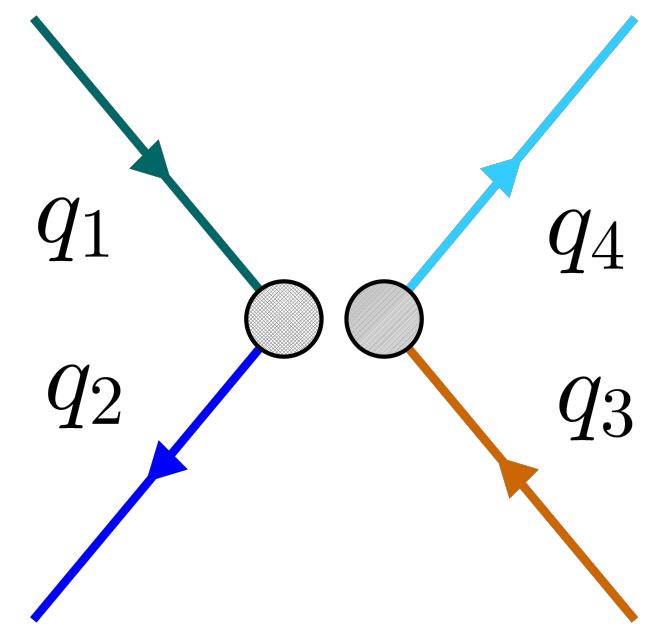


Gradient Flow

Example

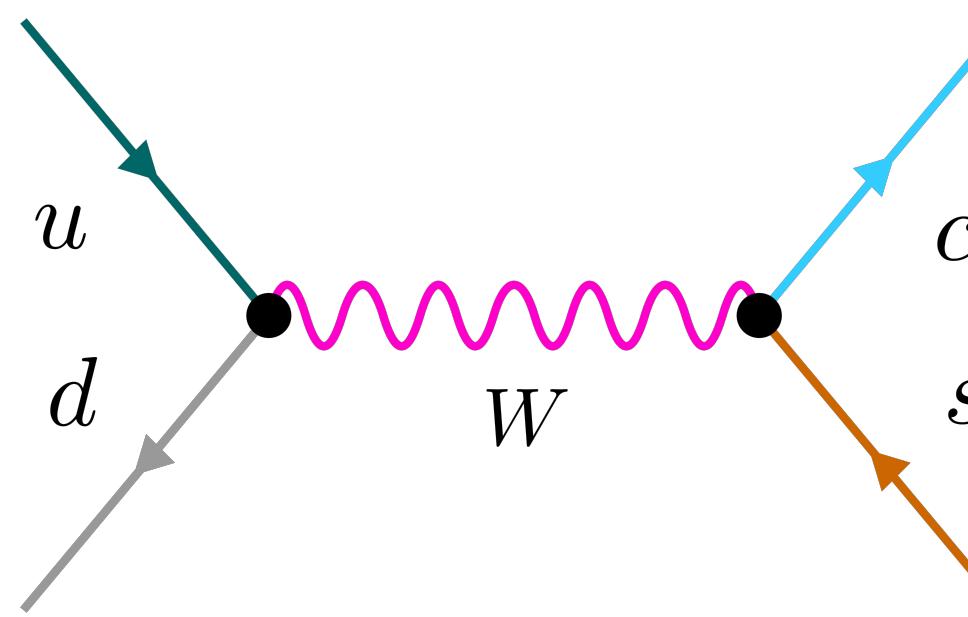


$$M_W \rightarrow \infty$$

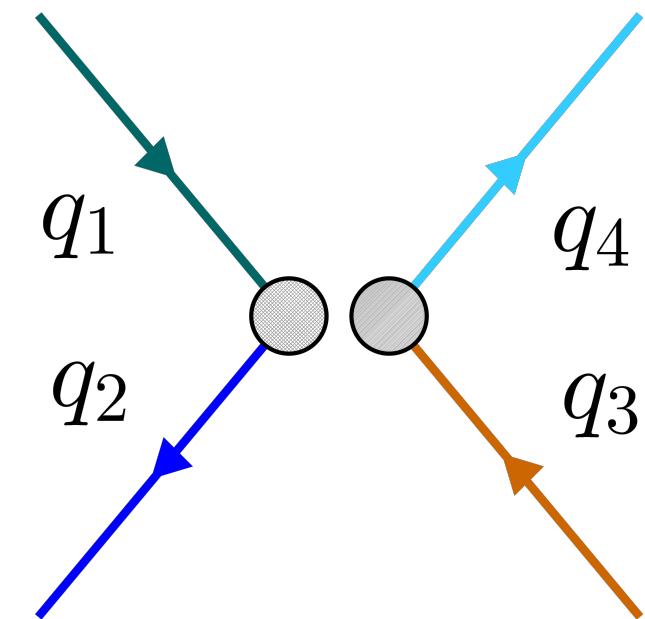


$$\begin{aligned}\mathcal{L}_{\text{eff}} &\ni \sum_n C_n^B \mathcal{O}_n \\ \langle T \rangle &= \sum_n C_n \langle \mathcal{O}_n^R \rangle\end{aligned}$$

Example



$$M_W \rightarrow \infty$$

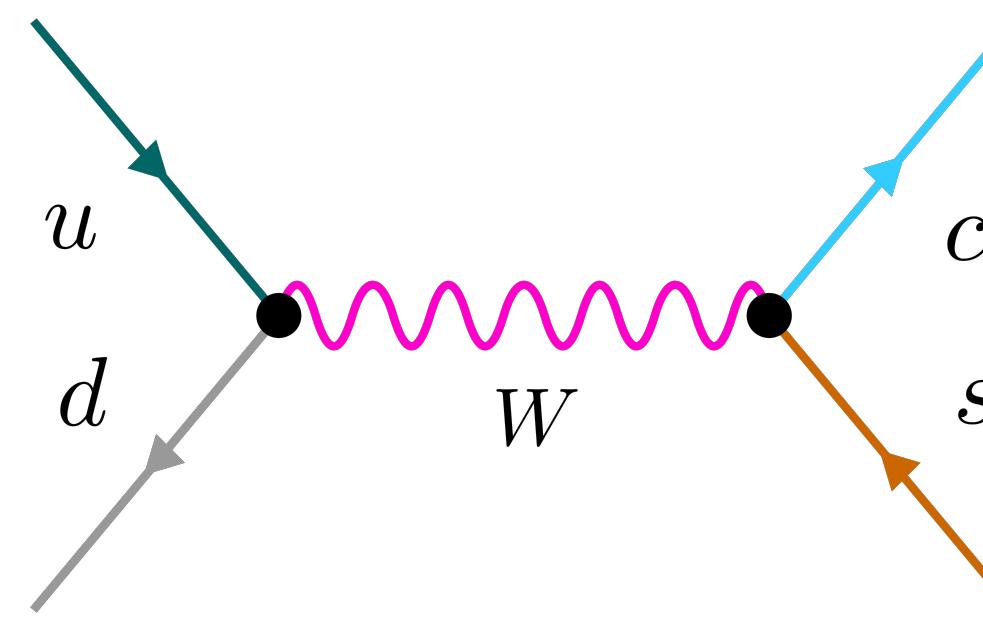



$$\mathcal{L}_{\text{eff}} \ni \sum_n C_n^B \mathcal{O}_n \equiv \sum_n \tilde{C}(\textcolor{red}{t})_n \tilde{\mathcal{O}}(\textcolor{red}{t})_n$$

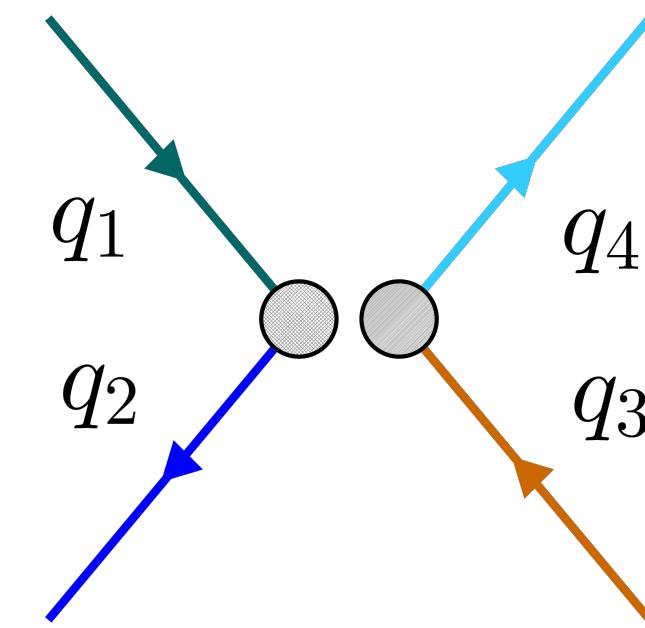
$$\langle T \rangle = \sum_n c_n \langle \mathcal{O}_n^R \rangle$$

$$\begin{aligned}\mathcal{O}_1 &= (\bar{q}_1 \gamma_\mu^L T q_2)(\bar{q}_3 \gamma_L^\mu T q_4) \\ \mathcal{O}_2 &= (\bar{q}_1 \gamma_\mu^L q_2)(\bar{q}_3 \gamma_L^\mu q_4)\end{aligned}$$

Example



$M_W \rightarrow \infty$



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pert.th. lattice

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\downarrow

$$\tilde{\mathcal{O}}_1(\textcolor{red}{t}), \tilde{\mathcal{O}}_2(\textcolor{red}{t})$$

Small-flow-time expansion

Lüscher, Weisz 2011

$$\tilde{\mathcal{O}}_n(\textcolor{red}{t}) \rightarrow \sum_m \zeta_{nm}^B(\textcolor{red}{t}) \mathcal{O}_m$$

perturbative

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Lüscher, Weisz 2011

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$\overline{\text{MS}}$ scheme

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Lüscher, Weisz 2011

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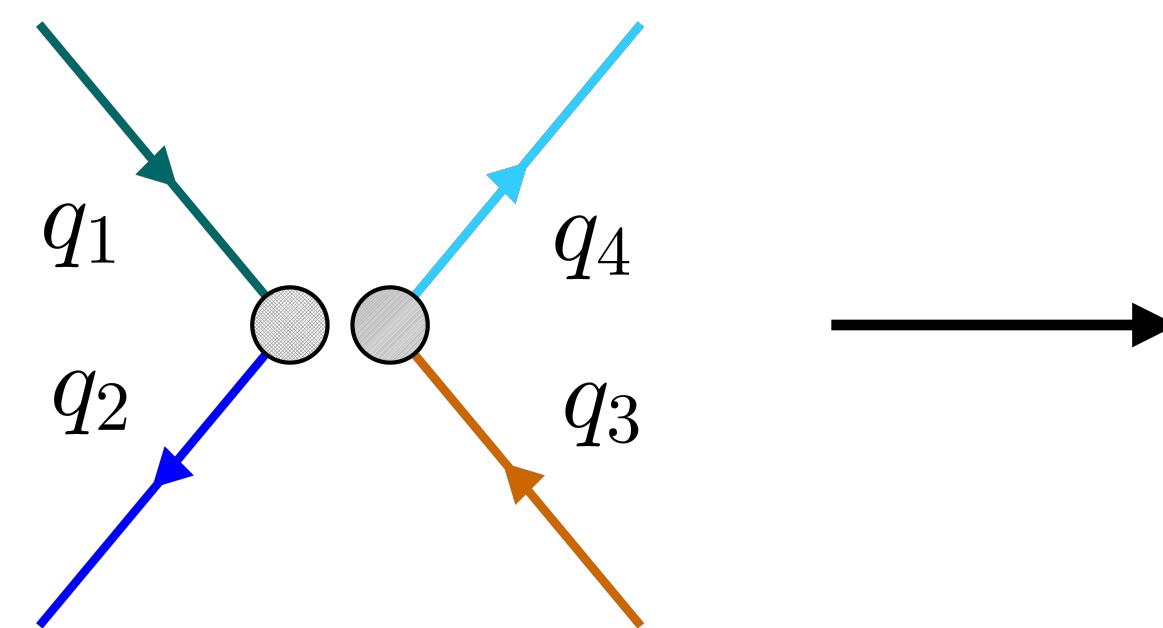
$$\tilde{C}_n(\textcolor{red}{t}) = C_m \zeta_{mn}^{-1}(\textcolor{red}{t})$$

Need to compute suitable Green's functions of the operators...

Method of projectors

$$\tilde{\mathcal{O}}_n(\textcolor{red}{t}) \rightarrow \sum_m \zeta_{nm}^B(\textcolor{red}{t}) \mathcal{O}_m \quad \longrightarrow \quad \langle k | \tilde{\mathcal{O}}_n(\textcolor{red}{t}) | 0 \rangle = \sum_m \zeta_{nm}^B(\textcolor{red}{t}) \langle k | \mathcal{O}_m | 0 \rangle$$

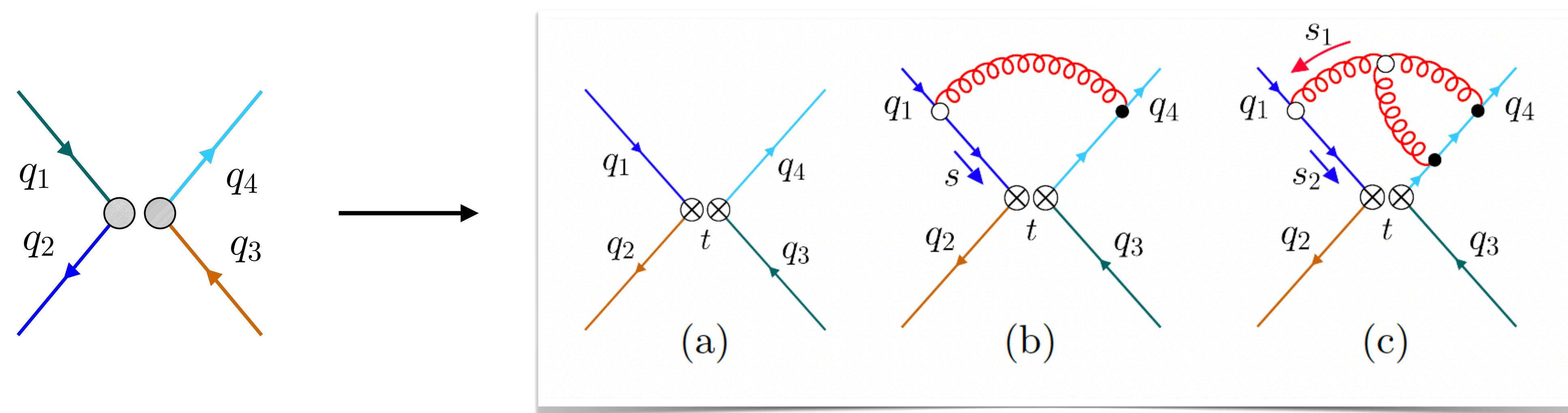
ideally: $\langle k | \mathcal{O}_m | 0 \rangle = \delta_{km} \Rightarrow \zeta_{nk}^B(\textcolor{red}{t}) = \langle k | \tilde{\mathcal{O}}_m(\textcolor{red}{t}) | 0 \rangle$



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Energy-momentum tensor

$$T_{\mu\nu}(x) = \frac{2}{|g(x)|^{1/2}} \frac{\delta S}{\delta g^{\mu\nu}(x)}$$

here: $S = S_{\text{QCD}}$ and $g^{\mu\nu}$ = flat metric

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Noether current of space-time translations

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Noether current of space-time translations
→ ill-defined on the lattice!

Energy-momentum tensor

in QCD:

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Energy-momentum tensor

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Energy-momentum tensor

in QCD:

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idea and NLO result: H. Suzuki '14

Flowed operators

$$\tilde{\mathcal{O}}_n(\textcolor{red}{t}) = \sum_m \zeta_{nm}(\textcolor{red}{t}) \mathcal{O}_m$$

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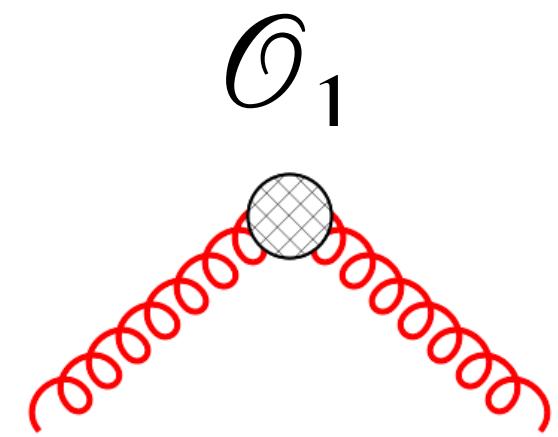
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e.g.

$$\mathcal{O}_{1,\mu\nu} = F_{\mu\rho}^a F_{\nu\rho}^a$$



$$-g^2 \delta^{ab} (\delta_{\mu\nu} p \cdot q - p_\mu q_\nu)$$

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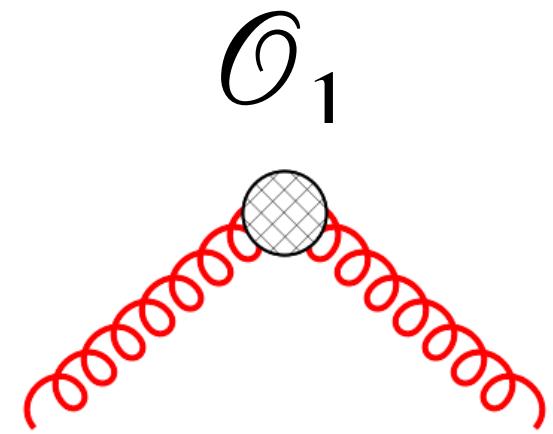
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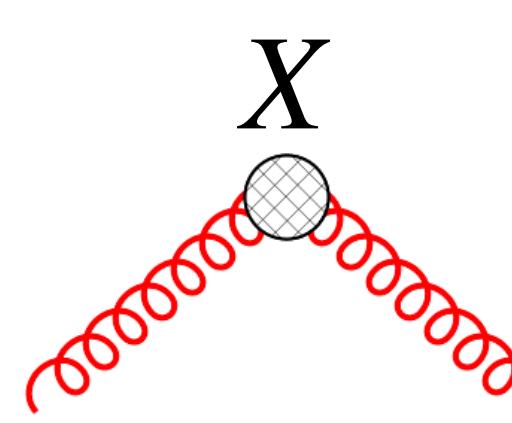
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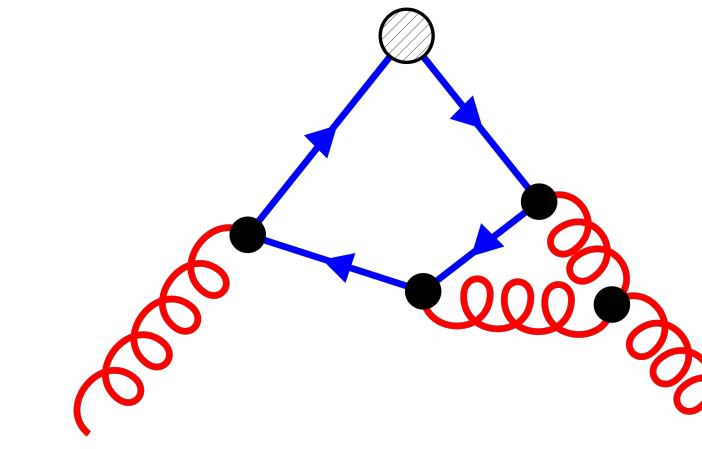
projector:

$$P_{1,\mu\nu}^{ab}[X] \sim \delta^{ab} \frac{\partial}{\partial p_\mu} \frac{\partial}{\partial q_\nu}$$



$$| p = q = 0$$

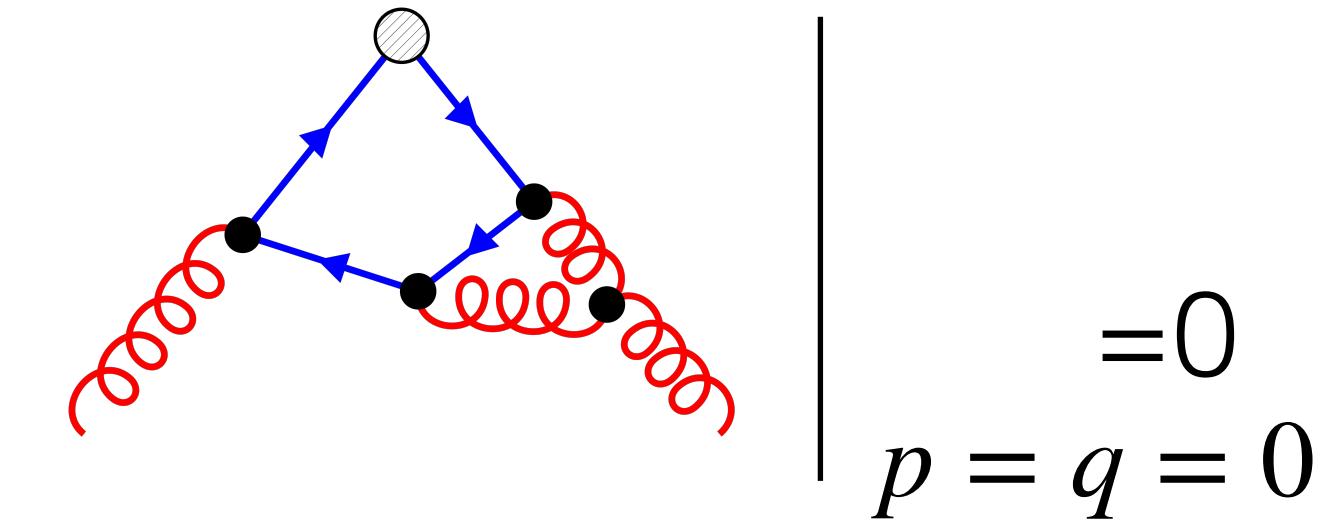
Method of projectors



$$\begin{array}{c} =0 \\ | \\ p = q = 0 \end{array}$$

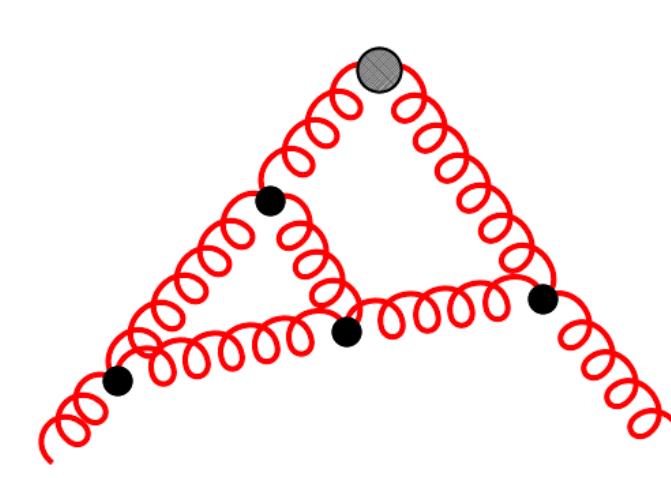
Method of projectors

If X is a regular QCD operator, all higher orders = 0.

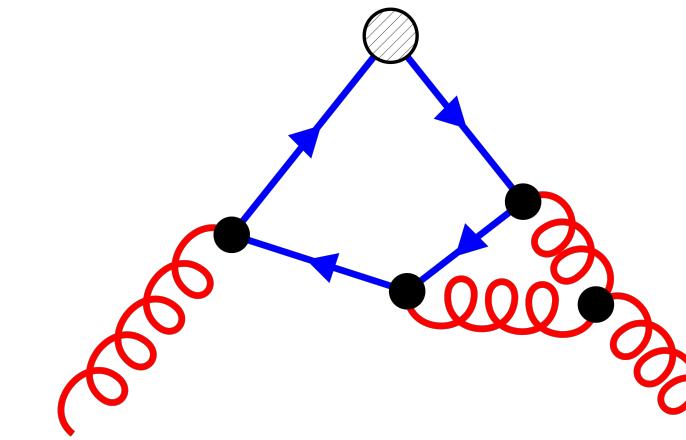


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$$\Rightarrow P_1[\tilde{\mathcal{O}}_n(\textcolor{red}{t})] = \sum_m \zeta_{nm}(\textcolor{red}{t}) P_1[\mathcal{O}_m] = \zeta_{n1}(\textcolor{red}{t})$$

to all orders

Gorishny, Larin, Tkachov '83

Energy-momentum tensor

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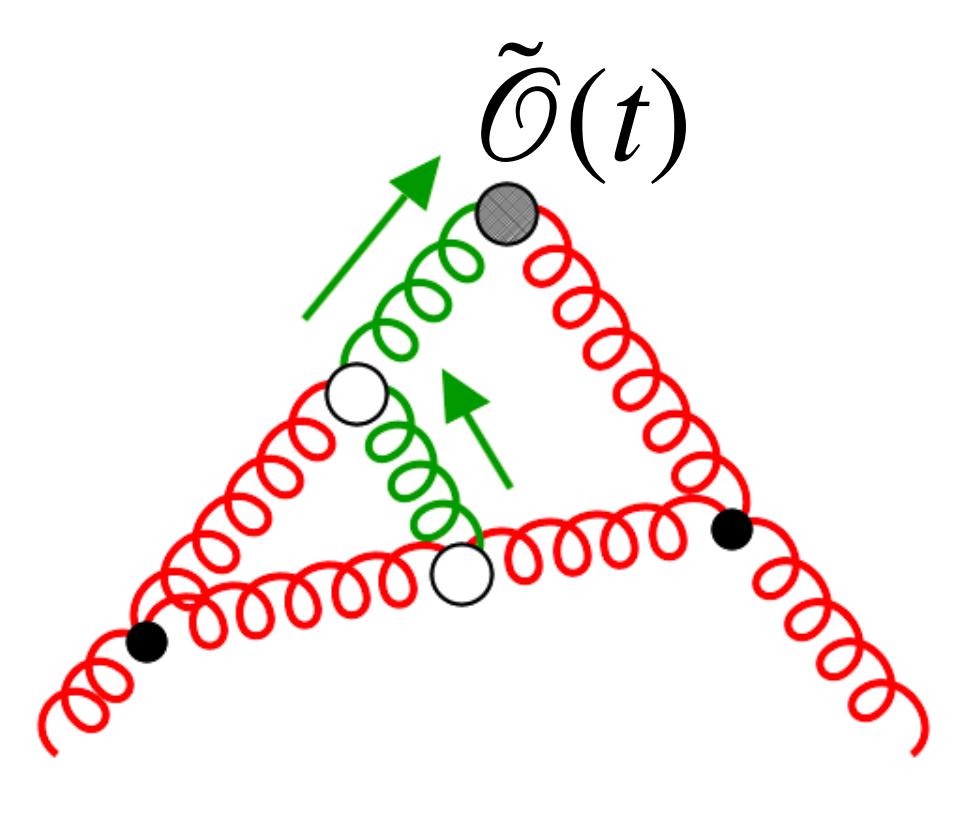
idea and NLO result: H. Suzuki '14

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$$\begin{aligned}\tilde{\mathcal{O}}_{1,\mu\nu}(t) &= F_{\mu\rho}^a(t) F_{\nu\rho}^a(t) \\ \tilde{\mathcal{O}}_{2,\mu\nu}(t) &= \delta_{\mu\nu} F_{\rho\sigma}^a(t) F_{\rho\sigma}^a(t) \\ \tilde{\mathcal{O}}_{3,\mu\nu}(t) &= \bar{\psi}(t) \left(\gamma_\mu \overleftrightarrow{D}_\nu(t) + \gamma_\nu \overleftrightarrow{D}_\mu(t) \right) \psi(t) \\ \tilde{\mathcal{O}}_{4,\mu\nu}(t) &= \delta_{\mu\nu} \bar{\psi}(t) \overleftrightarrow{D}(t) \psi(t)\end{aligned}$$

$$T_{\mu\nu}(x) = \sum_{n=1}^4 c_n(t) \tilde{\mathcal{O}}_{n,\mu\nu}(t, x)$$

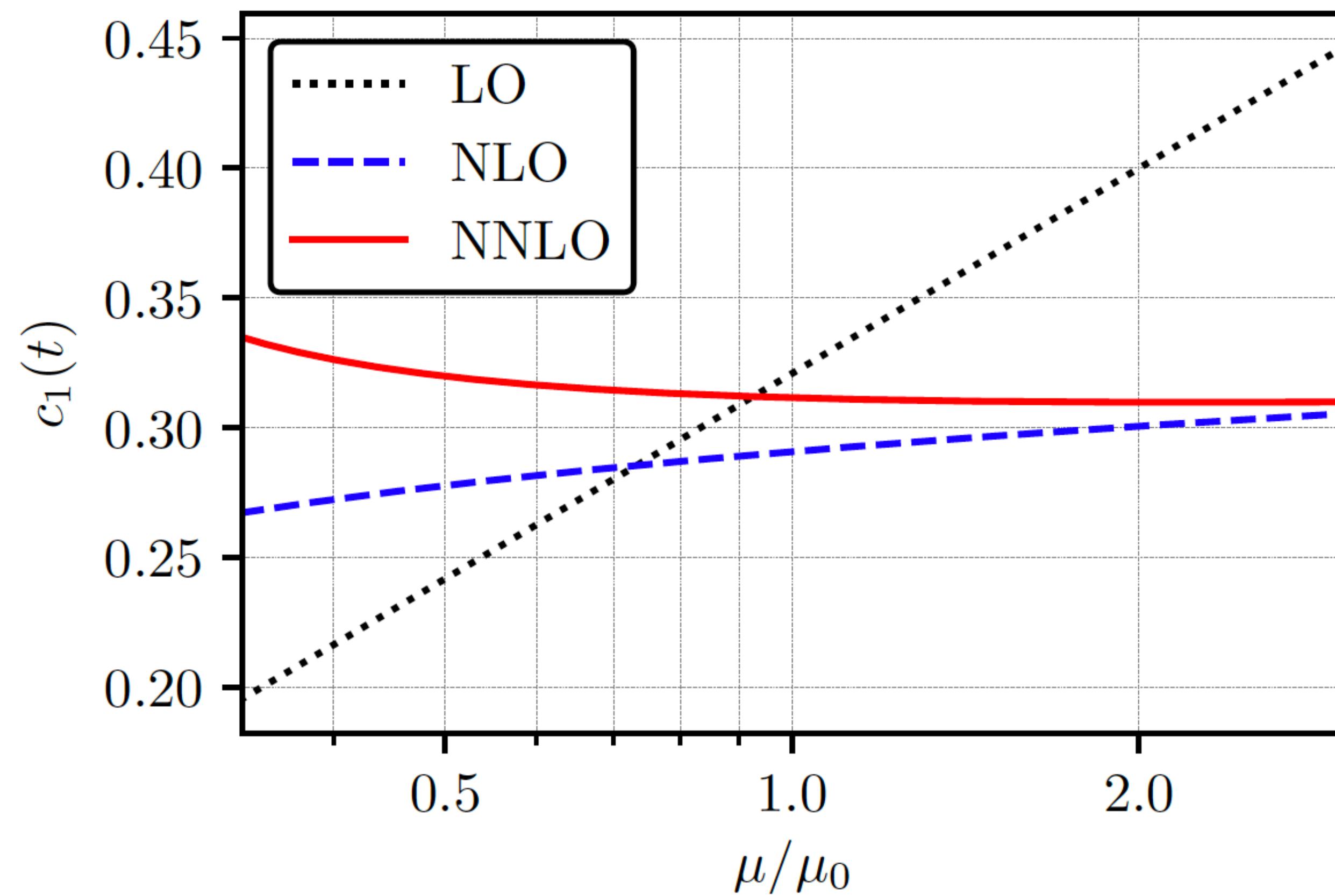
idea and NLO result: H. Suzuki '14

NNLO result

$$c_1(t) = \frac{1}{g^2} \left\{ 1 + \frac{g^2}{(4\pi)^2} \left[-\frac{7}{3}C_A + \frac{3}{2}T_F - \beta_0 L(\mu, t) \right] \right.$$
$$+ \frac{g^4}{(4\pi)^4} \left[-\beta_1 L(\mu, t) + C_A^2 \left(-\frac{14482}{405} - \frac{16546}{135} \ln 2 + \frac{1187}{10} \ln 3 \right) \right.$$
$$+ C_A T_F \left(\frac{59}{9} \text{Li}_2 \left(\frac{1}{4} \right) + \frac{10873}{810} + \frac{73}{54} \pi^2 - \frac{2773}{135} \ln 2 + \frac{302}{45} \ln 3 \right)$$
$$+ C_F T_F \left(-\frac{256}{9} \text{Li}_2 \left(\frac{1}{4} \right) + \frac{2587}{108} - \frac{7}{9} \pi^2 - \frac{106}{9} \ln 2 - \frac{161}{18} \ln 3 \right) \left. \right]$$
$$\left. + \mathcal{O}(g^6) \right\}, \quad L(\mu, t) \equiv \ln(2\mu^2 t) + \gamma_E$$

etc.

RH, Kluth, Lange '18

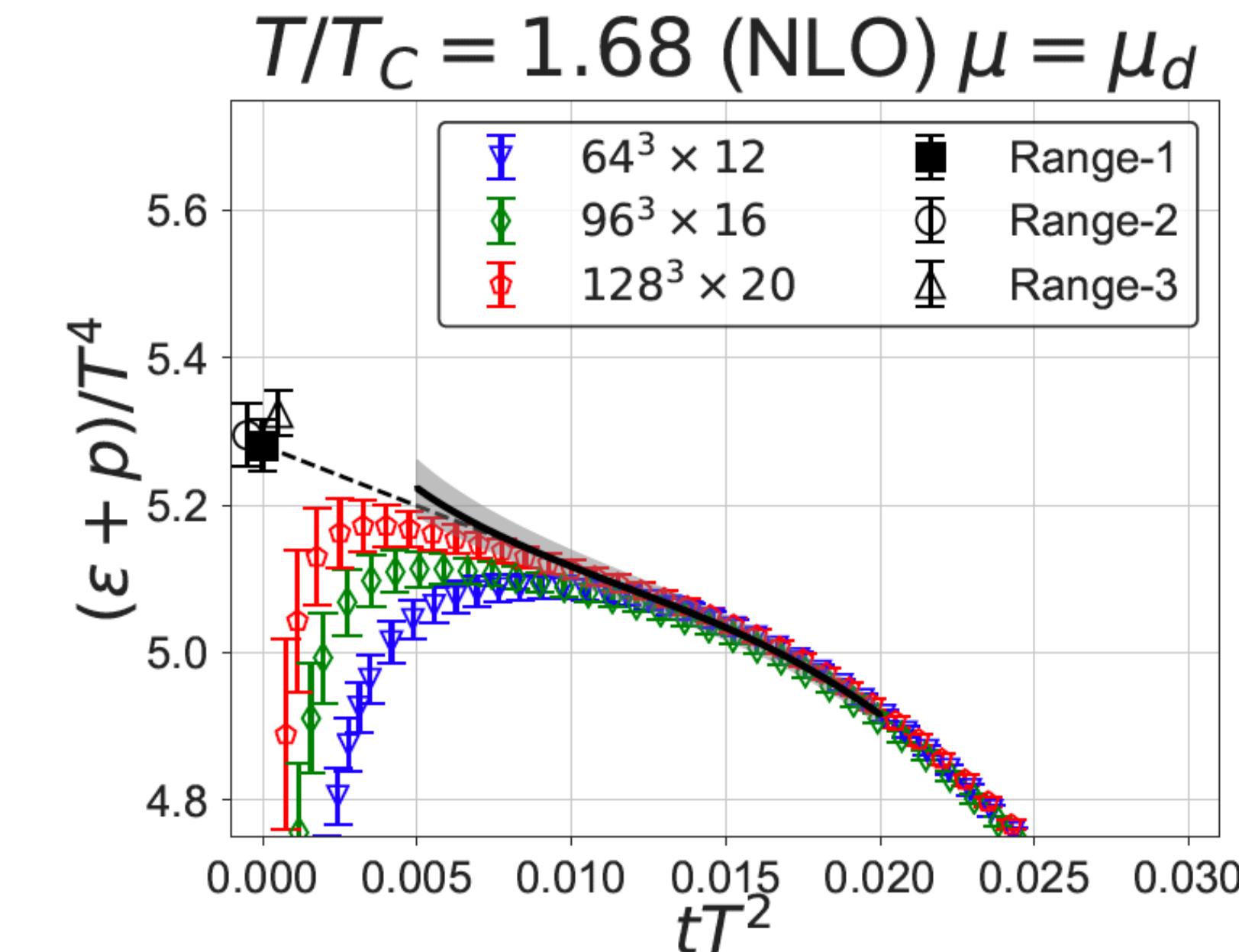
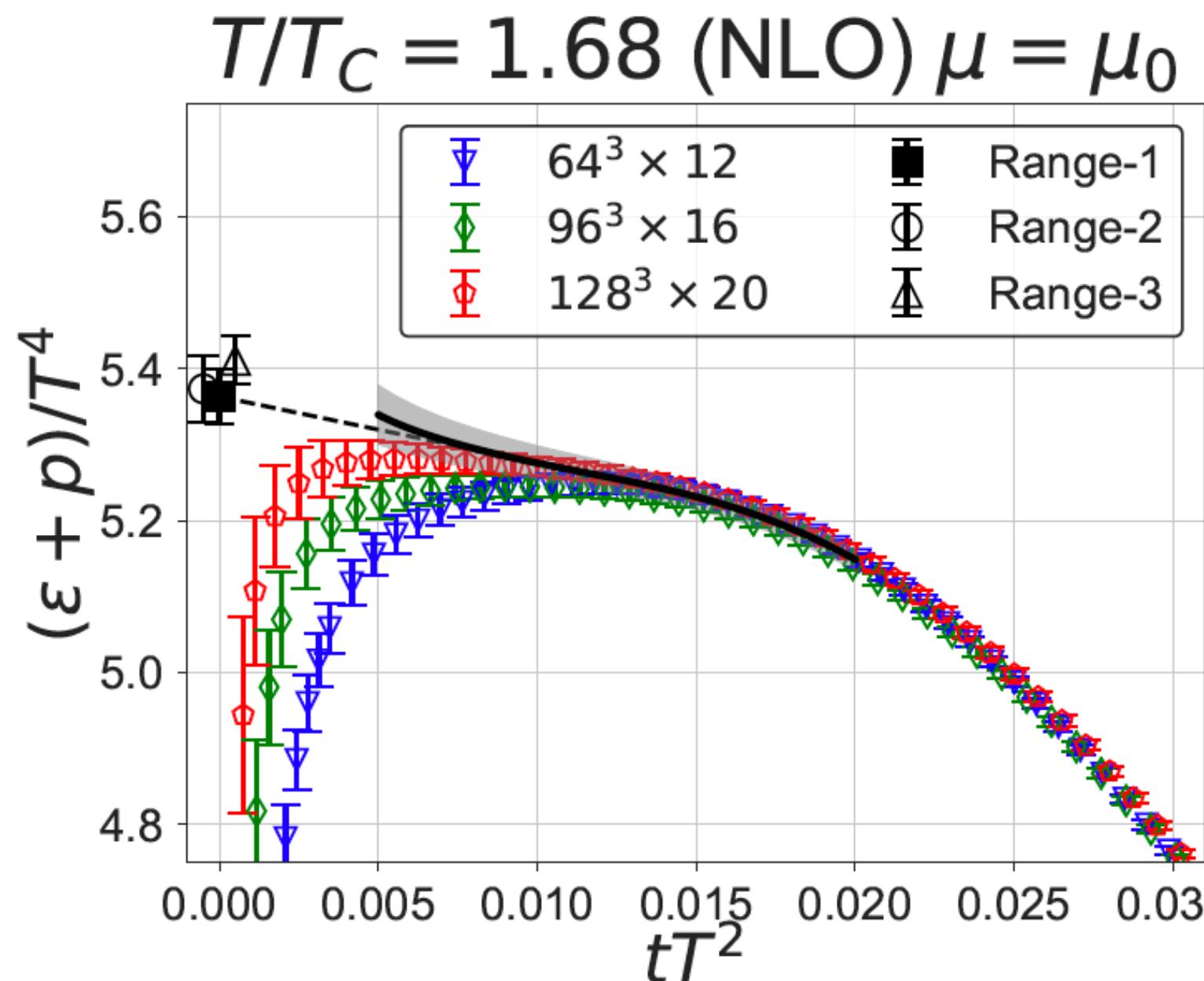
$\mu_0 = 3 \text{ GeV}$ 

Application

Entropy density: $\varepsilon + p = -\frac{4}{3} \left\langle T_{00}(x) - \frac{1}{4} T_{\mu\mu}(x) \right\rangle$

$$\mu_0 = \frac{e^{-\gamma_E/2}}{\sqrt{2t}}$$

$$\mu_d = \frac{1}{\sqrt{8t}}$$

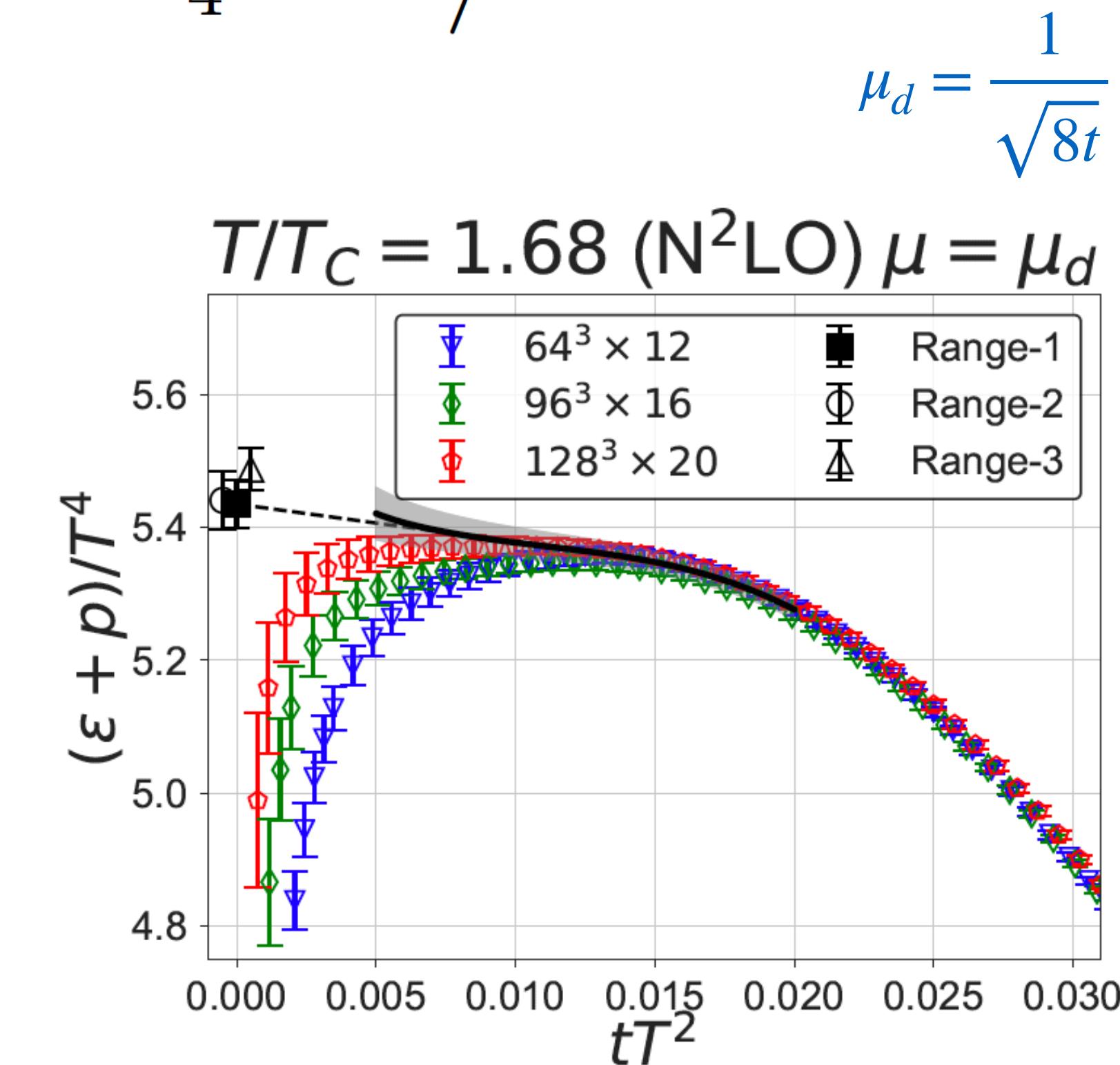
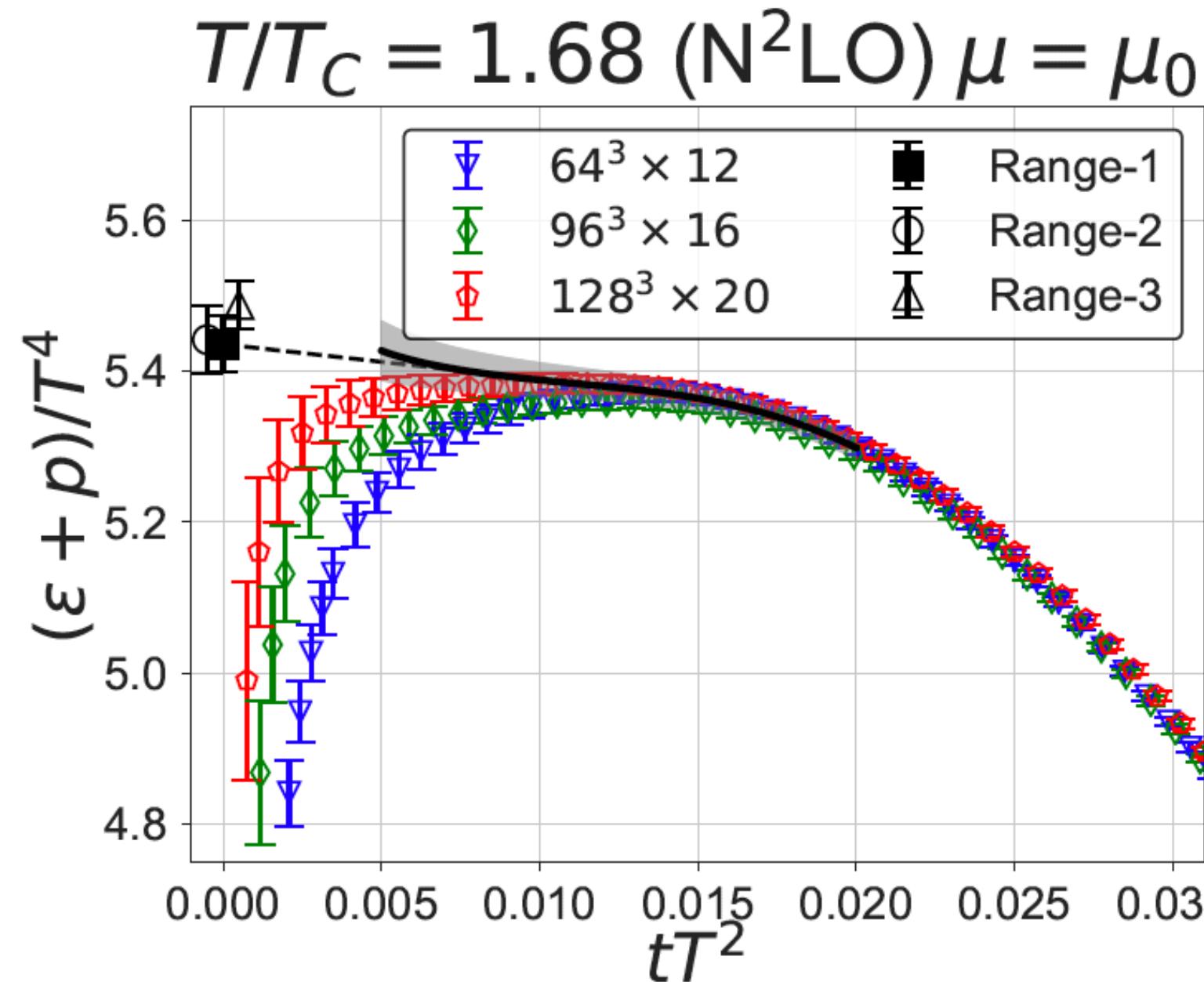


Iritani, Kitazawa, Suzuki, Takaura 2019

Application

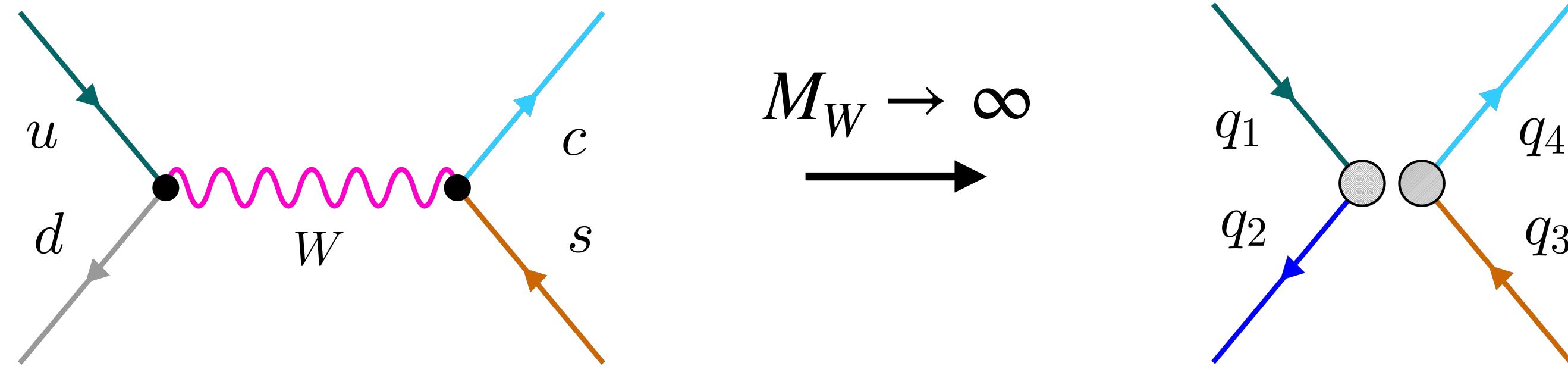
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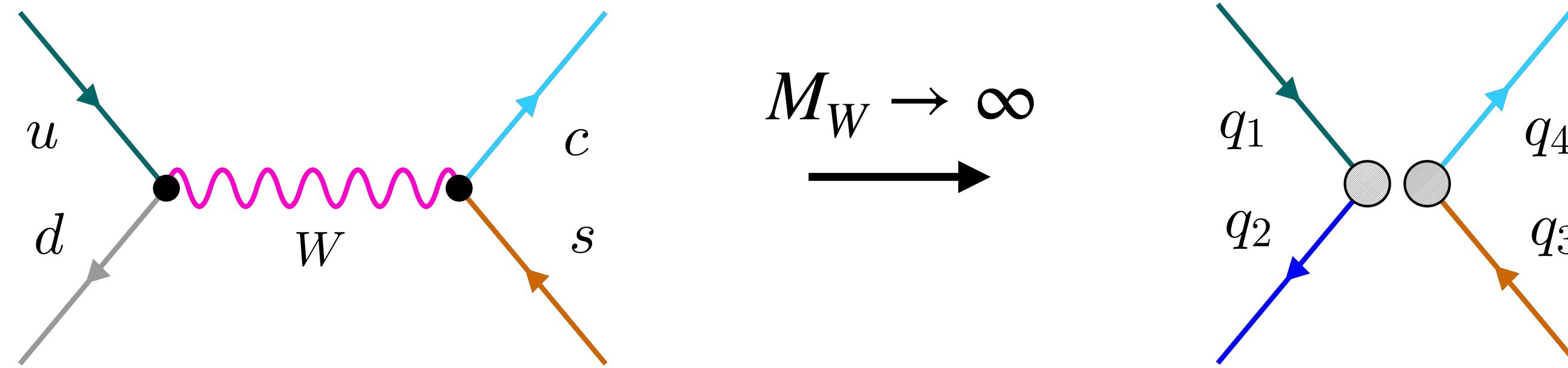
Iritani, Kitazawa, Suzuki, Takaura 2019

Application to EFT



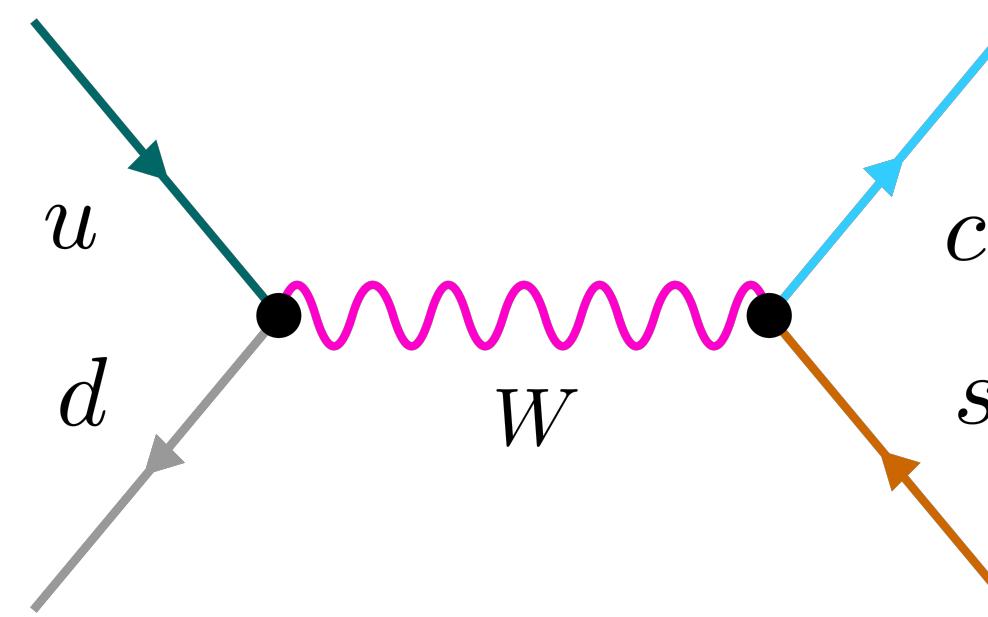
$$\sum_n C_n^B \mathcal{O}_n$$

Application to EFT

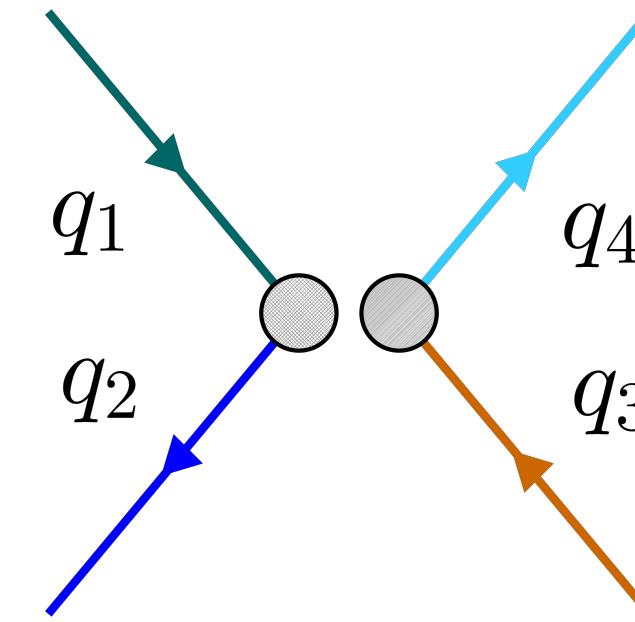


$$\sum_n C_n^B \mathcal{O}_n \equiv \sum_n \tilde{C}(\textcolor{red}{t})_n \tilde{\mathcal{O}}(\textcolor{red}{t})_n$$

Application to EFT

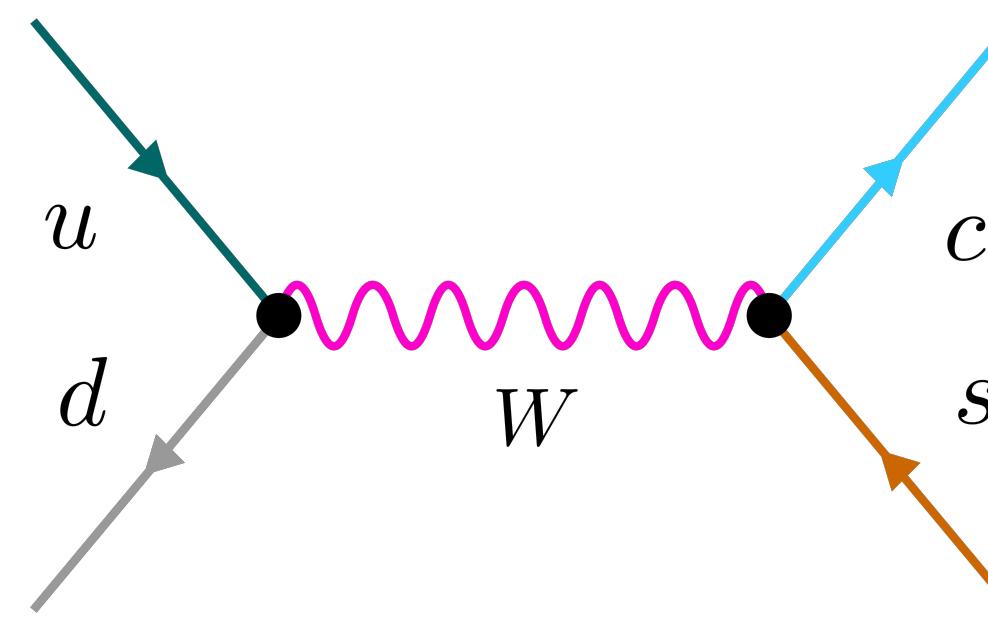


$$M_W \rightarrow \infty$$

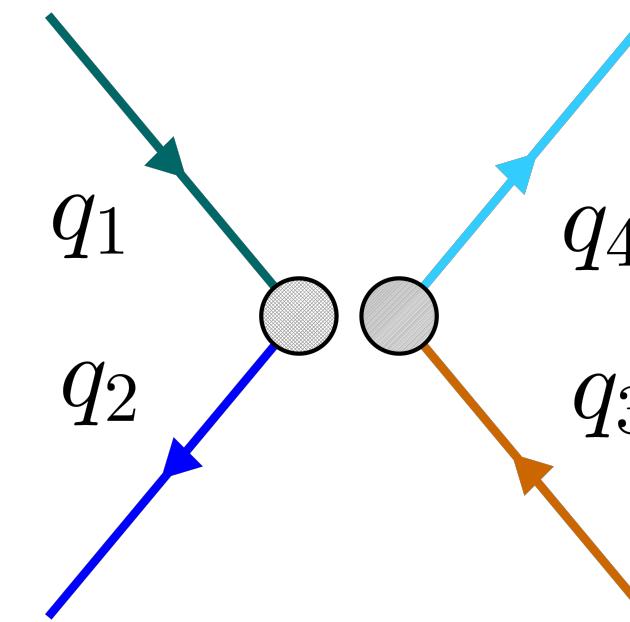


$$\sum_n C_n^B \mathcal{O}_n \equiv \sum_n \tilde{C}(\textcolor{red}{t})_n \tilde{\mathcal{O}}(\textcolor{red}{t})_n = \sum_n (C\zeta^{-1}(t))_n (\zeta(t)\mathcal{O}^R)_n$$

Application to EFT



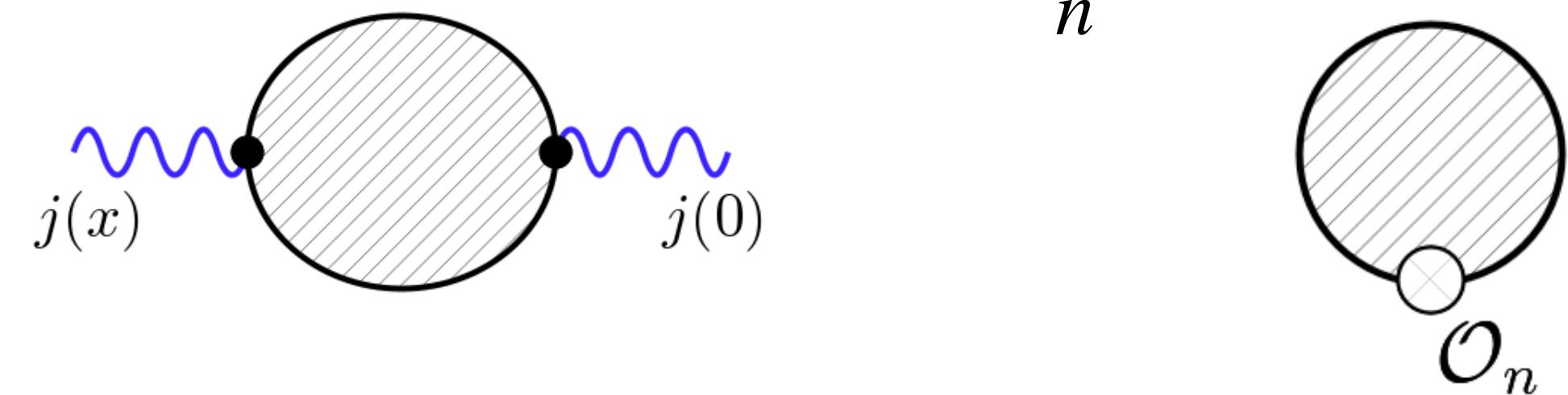
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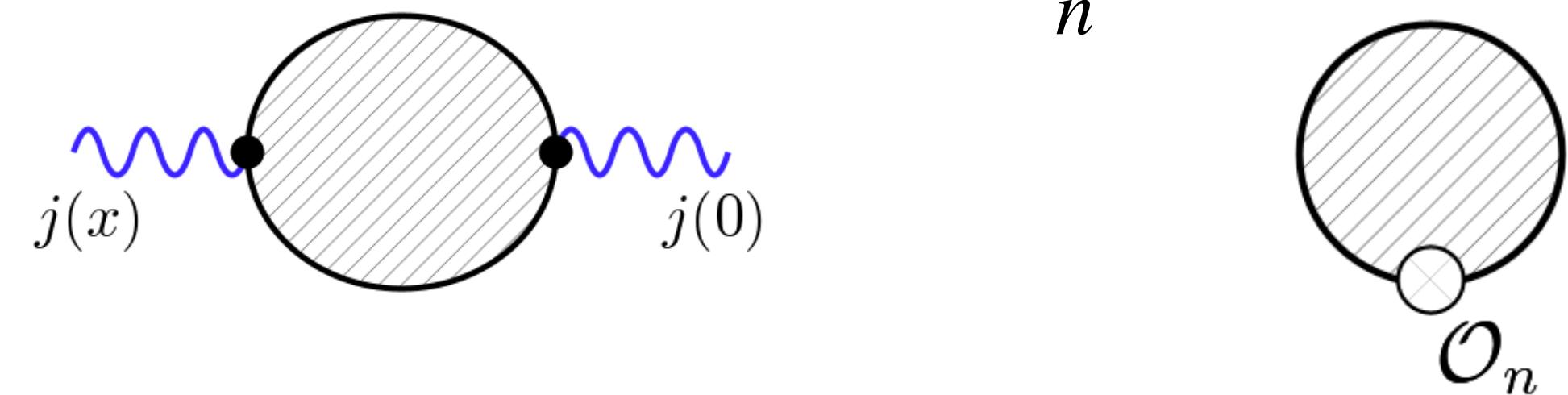
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$\overline{\text{MS}}$ coefficient

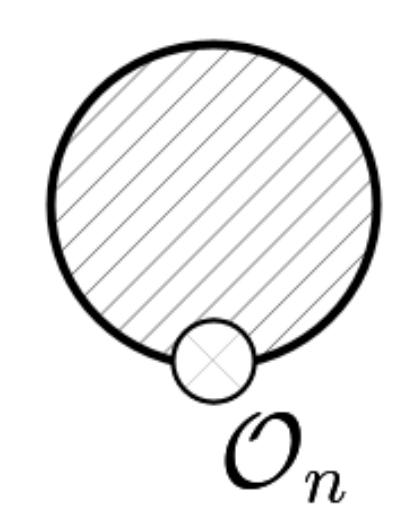
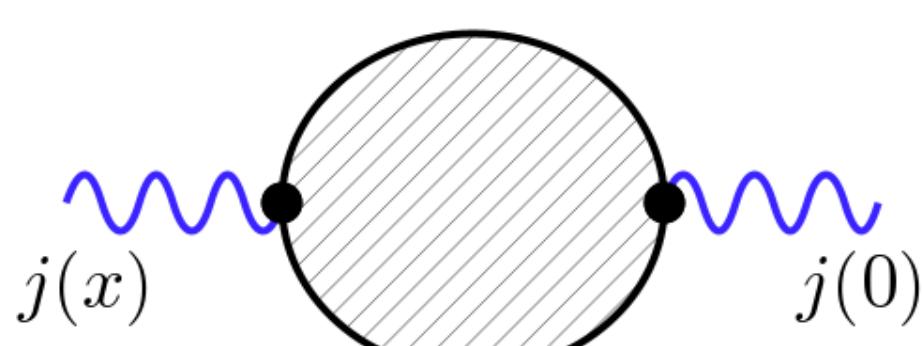
Hadronic vacuum polarization

$$\int d^4x e^{iQx} \langle T j(x) j(0) \rangle \rightarrow \sum_n C_n(Q) \langle \mathcal{O}_n(x=0) \rangle$$


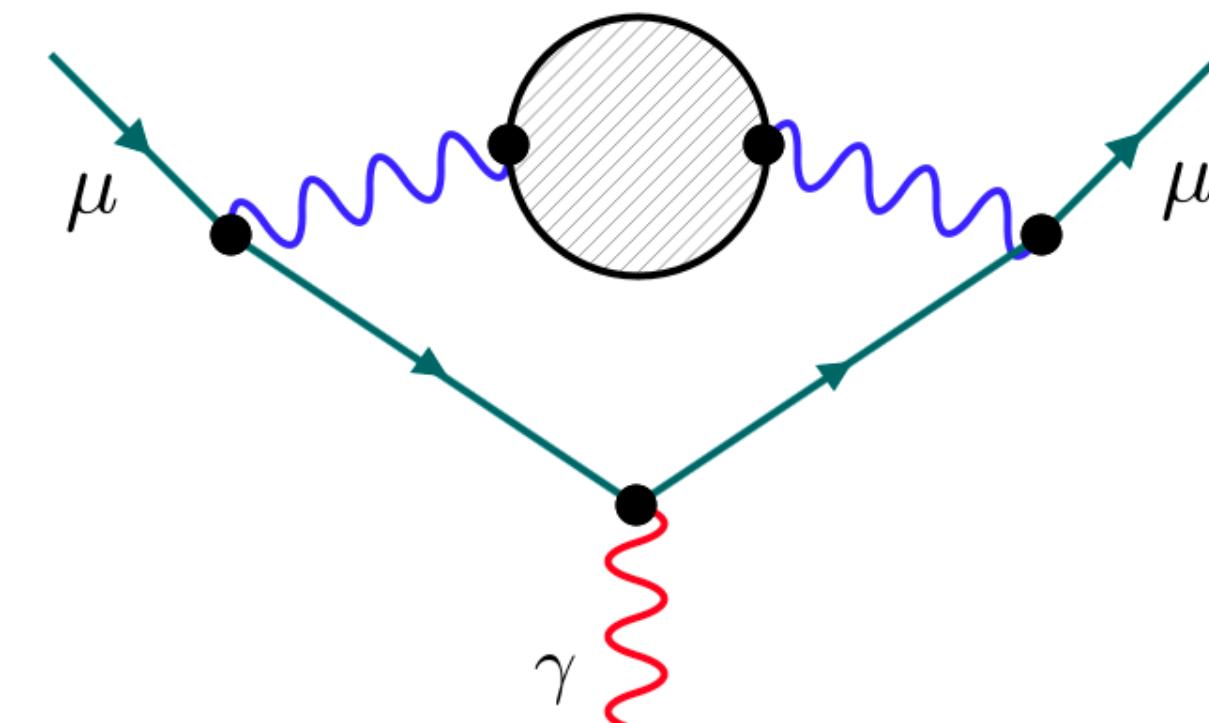
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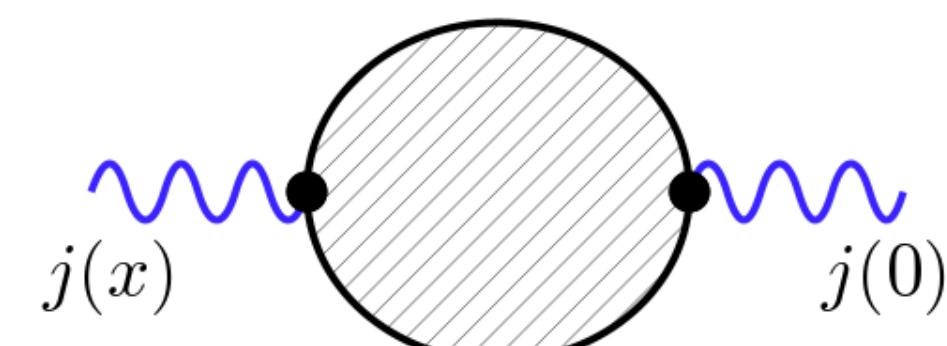
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contribution to $(g-2)\mu$

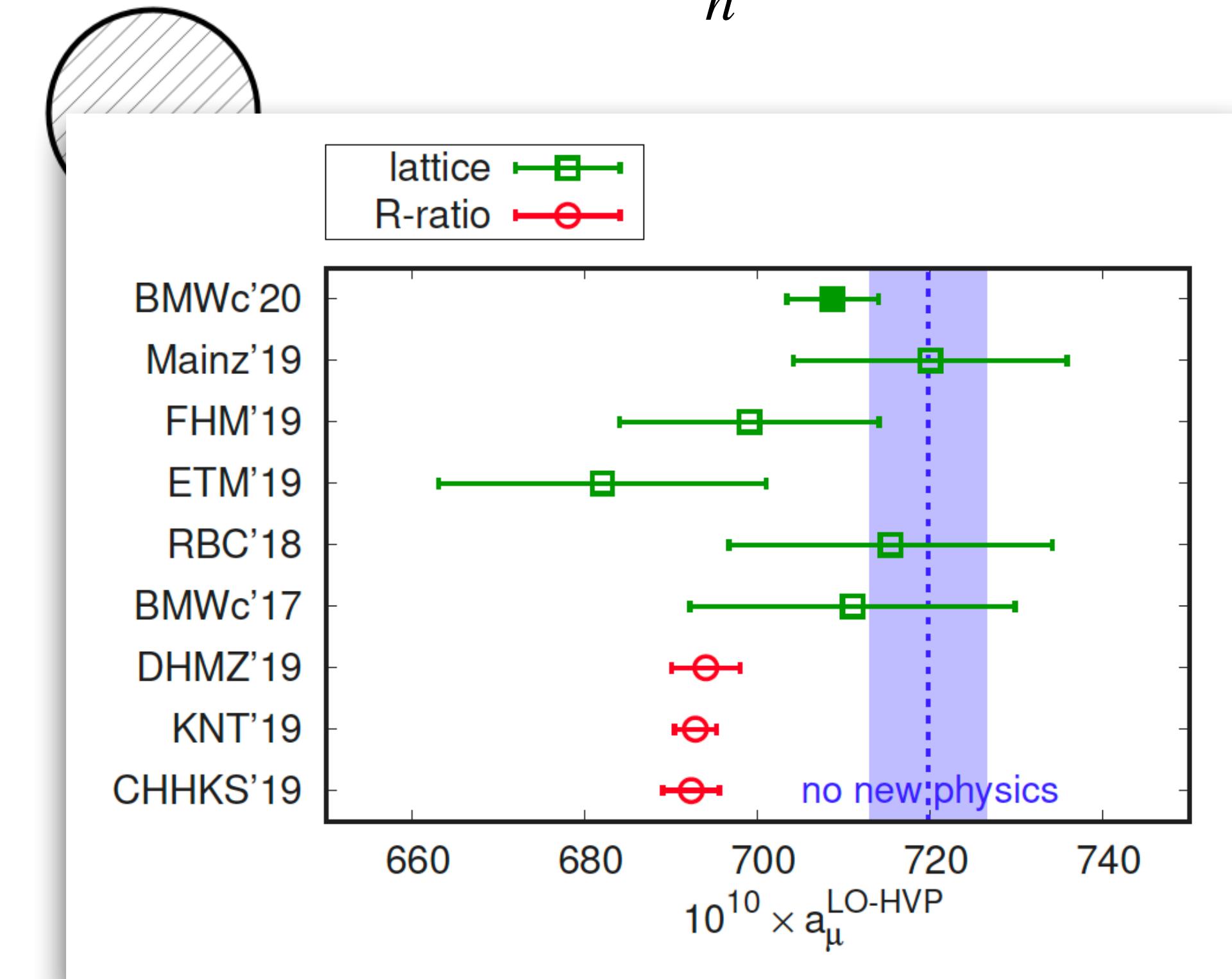
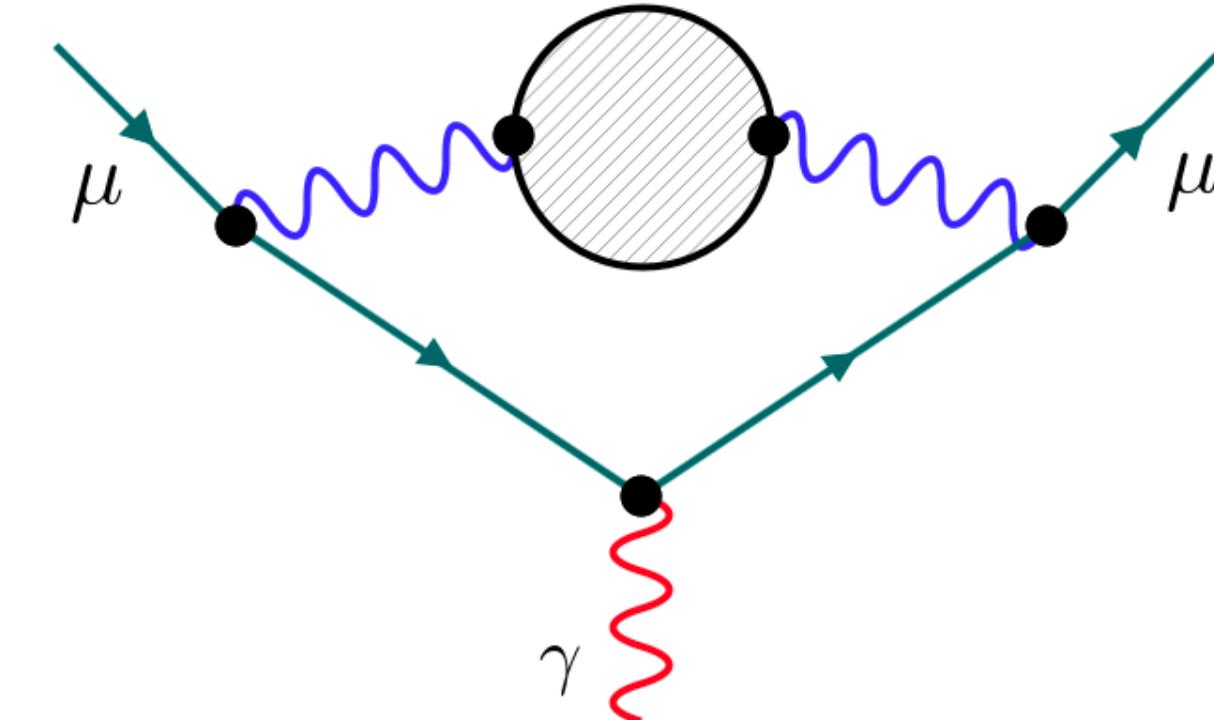


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contribution to $(g-2)\mu$



Renormalization

$$\tilde{\mathcal{O}}_n(\textcolor{red}{t}) = \sum_m \zeta_{nm}^B(\textcolor{red}{t}) \mathcal{O}_m$$

Renormalization

$$\tilde{\mathcal{O}}_n(\textcolor{red}{t}) = \sum_{\text{finite}} \zeta_{nm}^B(\textcolor{red}{t}) \mathcal{O}_m$$

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finite m divergent

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finite m divergent

$$\mathcal{O}^R = Z \mathcal{O}$$

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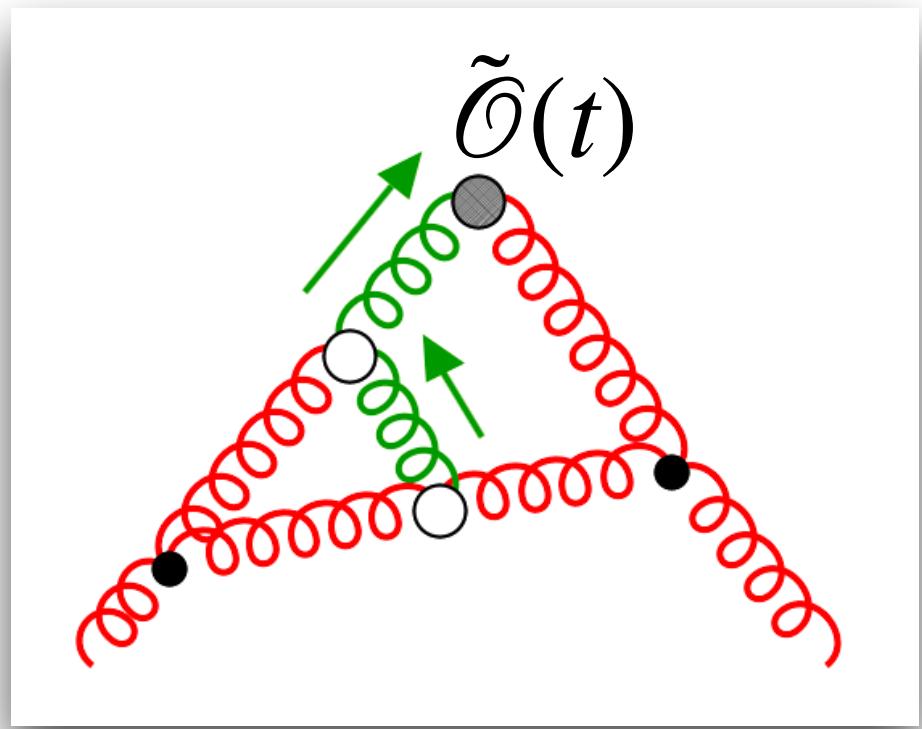
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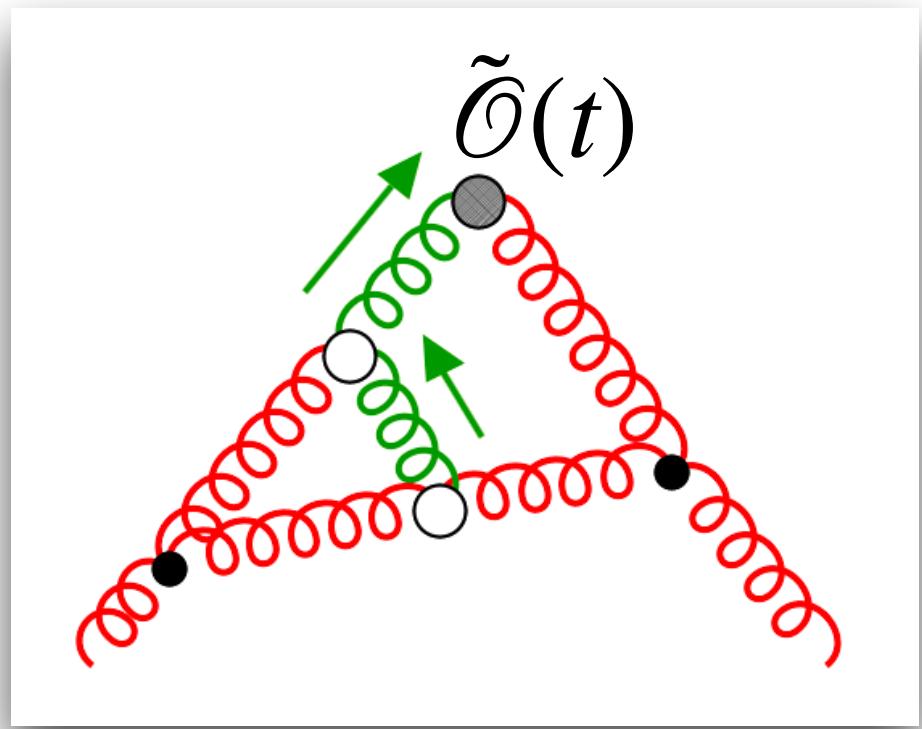
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$$\text{UV} = 0$$

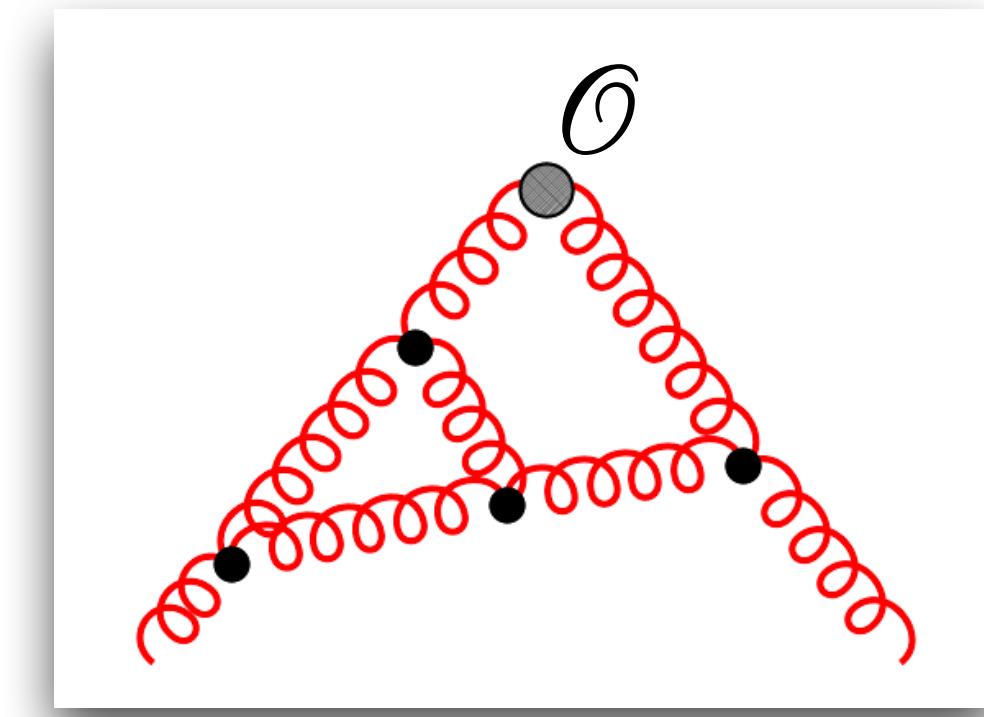
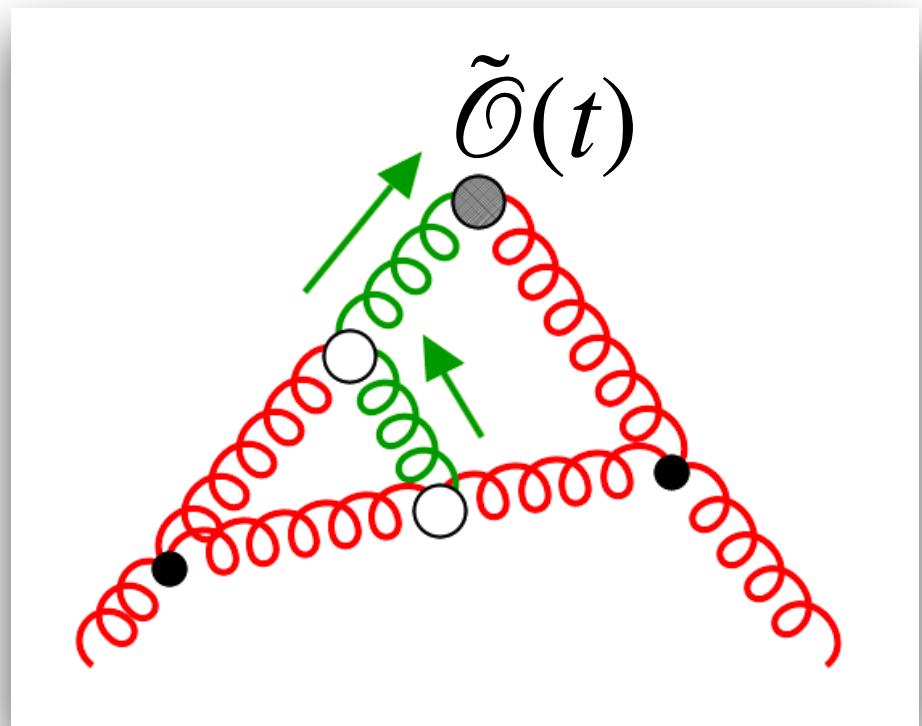
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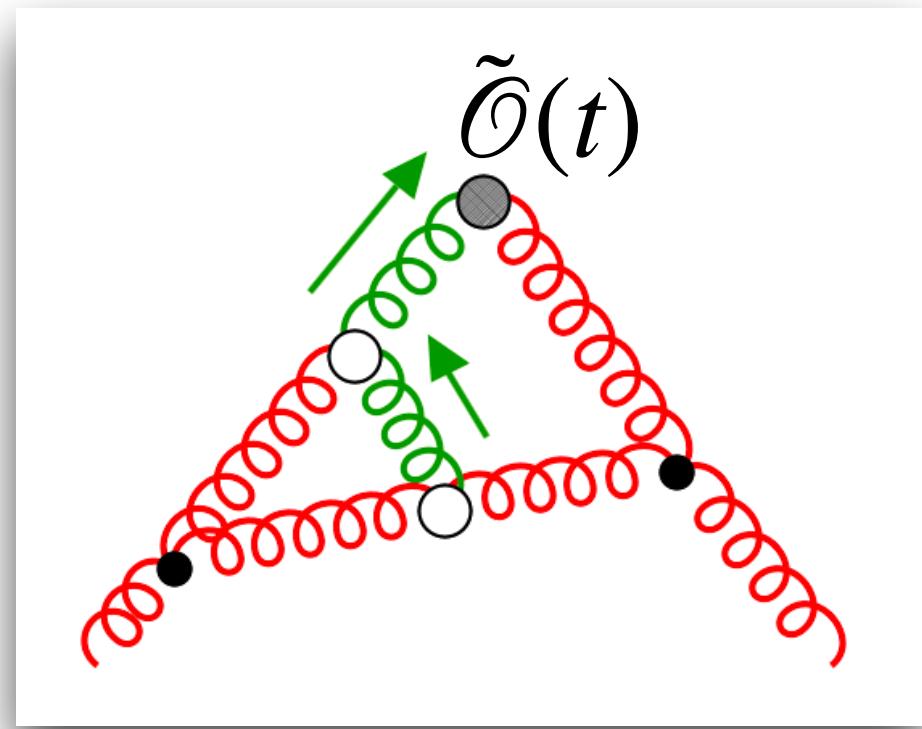
Renormalization

$$\tilde{\mathcal{O}}_n(\textcolor{red}{t}) = \sum_m \zeta_{nm}^B(\textcolor{red}{t}) \mathcal{O}_m$$

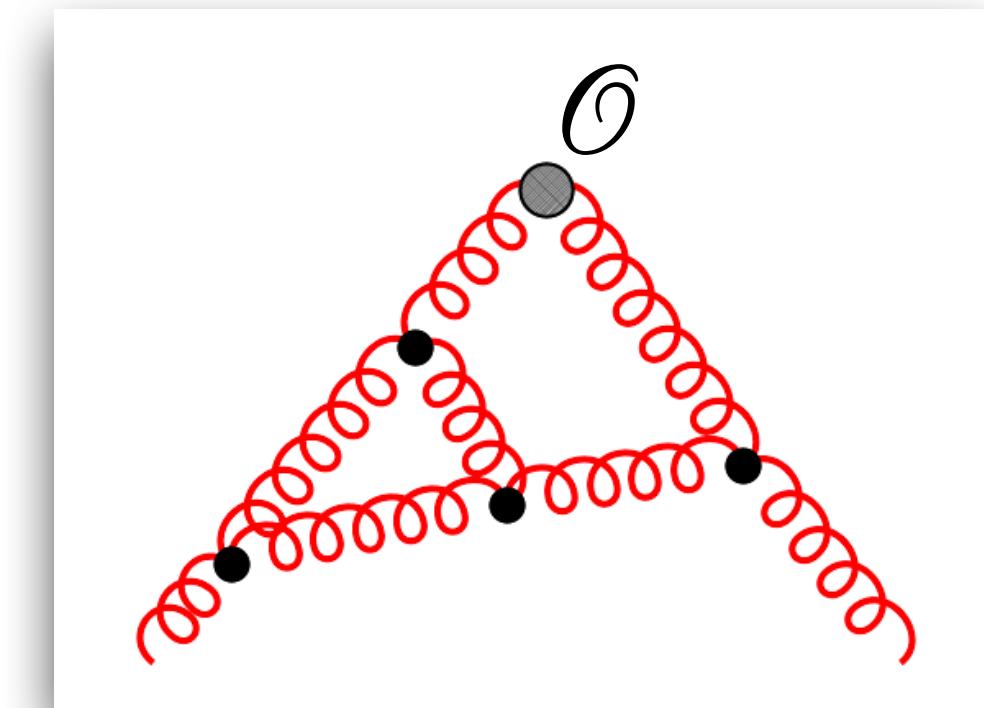
finite divergent

$$\mathcal{O}^R = Z \mathcal{O} \quad \tilde{\mathcal{O}}(\textcolor{red}{t}) = \zeta^B(\textcolor{red}{t}) Z^{-1} Z \mathcal{O} = \zeta^B(\textcolor{red}{t}) Z^{-1} \mathcal{O}^R = \zeta(\textcolor{red}{t}) \mathcal{O}^R$$

$$P_k[\tilde{\mathcal{O}}_n(\textcolor{red}{t})] = \sum_m \zeta_{nm}^B(\textcolor{red}{t}) P_k[\mathcal{O}_m] = \zeta_{nk}^B(\textcolor{red}{t})$$



UV = 0



UV + IR = 0

Summary

- Gradient Flow is a (relatively) new tool
- Extremely successful in lattice QCD
- Perturbative approach not yet fully explored

