#### Complexity + Simplicity in Theoretical Particle Physics

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supported by

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# The complexity and simplicity of Feynman diagrams

















 $\mathcal{L} = -\frac{1}{4} F_{r'} F^{\mu \nu}$  $+i\overline{\psi}\overline{D}\psi + h.c.$ ) + h.c.  $|\mathcal{D}_{\mu}\varphi|^{2} - V(\varphi)$ +





 $\mathcal{L} = -\frac{1}{4} F_{\mu} F^{\mu\nu}$  $+i\overline{\psi}\overline{D}\psi + h.c.$  $-\overline{\psi}\overline{\psi}\psi + h.c.$ +  $|\mathcal{D}_{\mu}\varphi|^2 - V(\varphi)$ 





#### Large Hadron Collider





## We are good...







Electron in magnetic field:  $E = - \overrightarrow{\mu} \cdot \overrightarrow{B}$ 







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What is Dark Matter? What is Dark Energy? Naturalness?





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e.g. Technicolor

New interaction instead of Higgs





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#### Theoretical Particle Physics is complex.





t=0: Particle at x with momentum p. *Question*: where is the particle at t > 0?





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101 (B) (2) (2) 2 OLO





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*Question*: where is the particle at t > 0?

$$\overrightarrow{F} = m\overrightarrow{a} = \frac{d}{dt}\overrightarrow{p}$$



















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$$S[path] = \int dt L(t)$$







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action:  $S[path] = \int dt L(t) \quad \Rightarrow \quad \overrightarrow{F} = m \overrightarrow{a}$ 




















probability amplitude

"Path integral"









Each path is weighted by exp(iS[path]).

$$S[\text{path}] = \int dt L(t) \in \mathbb{R}$$
$$L(t) = E_{kin}(t) - E_{pot}(t)$$

$$A(x,0;x',t) = \int_{P:x(0)\to x'(t)} dP \, e^{iS[P]} \in \mathbb{C}$$
  
"Path integra

probability amplitude

probability: 
$$|A(x,0;x',t)|^2$$

ſ





#### **Double slit experiment**







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source: <u>https://en.wikipedia.org/wiki/Double-slit\_experiment</u>





Recall classical mechanics:

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Recall classical mechanics:

t=0: Particle at x with momentum p. *Question*: where is the particle at t > 0?

Quantum Mechanics:

t=0: Particle at x with momentum p. *Question*: what is the probability to find particle at x' at t > 0?











#### 101 101 121 121 2 OLO





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Heisenberg uncertainty principle:  $(\Delta x)(\Delta p) > \hbar/2$ Classical state cannot be sharply defined!







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Even classically only probability can be calculated, but this is systematics!





# **Evolution of \psi(x)**

Evolve each x along all possible paths.







So far: one particle in one space dimension: 1 d.o.f. Things may change with coupled (interacting) d.o.f.'s:







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action:  
$$S[P_1,P_2] = S_1[P_1] + S_2[P_2] + S_{int}[P_1,P_2]$$





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action:  

$$S[P_{1},P_{2}] = S_{1}[P_{1}] + S_{2}[P_{2}] + S_{int}[P_{1},P_{2}]$$

$$\int_{P_{1}: x_{1}(0) \to x_{1}'(t)} \int_{P_{2}: x_{2}(0) \to x_{2}'(t)} dP_{1}dP_{2} e^{iS[P_{1},P_{2}]}$$





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But: principles remain the same (with a few additions...).



























Field theory: continuum limit  $\phi(x,0) \rightarrow \phi'(x,t)$ 









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 $\int \mathcal{D}\phi \, e^{iS[\phi]}$ 

how can we calculate that?





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Perturbation theory:  $S[\phi] = S_{\text{free}}[\phi] + S_{\text{int}}[\phi]$ 





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$$S_{\text{free}}[\phi] = \int_{x} \phi_x D_x \phi_x$$





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Perturbation theory:  $S[\phi] = S_{\text{free}}[\phi] + S_{\text{int}}[\phi]$ 

$$S_{\text{free}}[\phi] = \int_{x} \phi_{x} D_{x} \phi_{x} \qquad S_{\text{int}}[\phi] = \int_{x} g \phi_{x}^{3}$$
small





 $\int \mathcal{D}\phi \, e^{iS[\phi]}$  how can we calculate that?

Perturbation theory:  $S[\phi] = S_{\text{free}}[\phi] + S_{\text{int}}[\phi]$ 





#### free: no springs





 $\int \mathcal{D}\phi \, e^{iS[\phi]}$ 

how can we calculate that?











$$\int \mathscr{D}\phi \, e^{iS[\phi]} \quad \text{how can we calculate that?}$$
Using 
$$\int_{-\infty}^{\infty} dx \, e^{-ax^2} = \sqrt{\frac{\pi}{a}}$$

$$\int \mathscr{D}\phi \, e^{iS_{\text{free}}[\phi]} \quad \text{is calculable!}$$





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Result: 
$$\left(\sqrt{\frac{\pi}{D_x}}\right)$$





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$$\left(\sqrt{\frac{\pi}{D_x}}\right)^{\infty}$$




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Result: 
$$\left(\sqrt{\frac{\pi}{D_x}}\right)^{\infty\infty}$$





















































$$\int \mathscr{D}\phi \, e^{iS[\phi]} = \int \mathscr{D}\phi \, e^{iS_{\rm int}[\phi]} e^{iS_{\rm free}[\phi]}$$





$$\int \mathscr{D}\phi \, e^{iS[\phi]} = \int \mathscr{D}\phi \, e^{iS_{\text{int}}[\phi]} e^{iS_{\text{free}}[\phi]}$$
$$S_{\text{int}}[\phi] = g \int_{X} \phi_{X}^{3}$$





$$\int \mathscr{D}\phi \ e^{iS[\phi]} = \int \mathscr{D}\phi \ e^{iS_{\text{int}}[\phi]} \ e^{iS_{\text{free}}[\phi]}$$
$$S_{\text{int}}[\phi] = g \int_{x} \phi_{x}^{3}$$
$$\int \mathscr{D}\phi \left[ 1 + \phi_{x}^{3} + \frac{1}{2} \ \phi_{x}^{3} \phi_{y}^{3} + \frac{1}{3!} \ \phi_{x}^{3} \phi_{y}^{3} \phi_{z}^{3} + \dots \right] e^{iS_{\text{free}}[\phi]}$$



















## Scattering: external fields







Forget all of the above!





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Forget all of the above!



# Connect initial and final state with these objects $\Rightarrow$ probability amplitude





Forget all of the above!



Connect initial and final state with these objects  $\Rightarrow$  probability amplitude









 $-8\pi\alpha$ 





## **Z-boson production**







# **Higgs discovery**









# **Higgs discovery**











Forget all of the above!



Connect initial and final state with these objects  $\Rightarrow$  probability amplitude





























 $e^{\dagger}$ 





R. Harlander, Theoretical Particle Physics, May 2019





9

## **Standard Model**







## The virtues of Feynman diagrams



Easy to use.





## The virtues of Feynman diagrams



Easy to use.



Easy to talk about.





## The virtues of Feynman diagrams



Easy to use.



Easy to talk about.

Systematic and algorithmic:

theory: (g-2)/2 = 0.001159652181643(764)experiment: (g-2)/2 = 0.00115965218073(28)





#### We love them...











#### We love them...





sunrise/ sunset










sunrise/ sunset



#### watermellon











sunrise/ sunset



#### watermellon



tennis court









sunrise/ sunset



#### watermellon















penguin













#### Interpretation?





 $\mathcal{D}\phi e^{iS[\phi]}$ 





# Where the simplicity begins to end: higher orders





 $\alpha$ 



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e\_





























- many diagrams
- complicated loop integrals
- divergent integrals





प्रयंग पर्युंग पर्युंग पर्नुंग पर्युंग पर्देग पर्नेग بوقع بوقع بعرقه بوقه بوقع بوقع بوقع ഷ്ട്ര ഫ്ല ഫ്ല ഫ്ല ഫ്ല ഫ്ല وقها وقاح وقاح وحال وقا وقا وقا وقا പ്പം പ്പം എം എം എം എം എം. പ്പം പും എം എം എം എം എം 62 <u>(</u> 20 (m)B 200 76 94

5-loop contributions to the anomalous magnetic moment of the electron

Today's calculations: O(100.000) diagrams

Require high level of automation.

source: https://www.sciencedirect.com/science/article/pii/S0370269317305324











```
qgraf-3.1.4
output= 'qlist.out' ;
style= 'form.sty' ;
model= 'qcd';
in= fq,fQ ;
out= fq,fQ ;
loops= 1;
loop_momentum= ;
options = ;
 3P --- 2+ 1- --- 2N+ 1C-
 3V --- 3^3
         -----
 3^4 --- 130 diagrams
 total = 130 diagrams
```





```
qgraf-3.1.4
 output= 'qlist.out' ;
 style= 'form.sty' ;
model= 'qcd';
in= fq,fQ ;
out= fq,fQ ;
 loops= 2;
 loop_momentum= ;
options = ;
 3P --- 2+ 1- --- 2N+ 1C-
 3V --- 3^3
  3^6 --- 4058 diagrams
 total = 4058 diagrams
```











```
-#[ d222:
   -1
  *vx(fq(-1),fQ(-3),g(1))
  *vx(fq(-4),fQ(-2),g(2))
  *vx(fq(4),fQ(3),g(1))
  *vx(g(2),g(5),g(6))
  *vx(fq(3),fQ(7),g(5))
  *vx(fq(9),fQ(4),g(8))
  *vx(g(6),g(8),g(10))
  *vx(fq(7),fQ(9),g(10))
#] d222:
-#[ d223:
  *vx(fq(-1),fQ(-3),g(1))
  *vx(fq(-4),fQ(-2),g(2))
  *vx(g(1),g(3),g(4))
  *vx(fq(5),fQ(6),g(2))
  *vx(fq(7),fQ(5),g(3))
  *vx(fq(8),fQ(9),g(4))
  *vx(fq(6),fQ(8),g(10))
  *vx(fq(9),fQ(7),g(10))
-#] d223:
```

#### Who ever looks at 134146 Feynman diagrams?





## Surely you're joking...





